

spontaneous time-reversal symmetry breaking in the quantum spin Hall effect

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PRL 96, 106401 (2006)

PHYSICAL REVIEW LETTERS

also Xu & Moore PRB 73, 045322 (2006)

Helical Liquid and the Edge of Quantum Spin Hall Systems

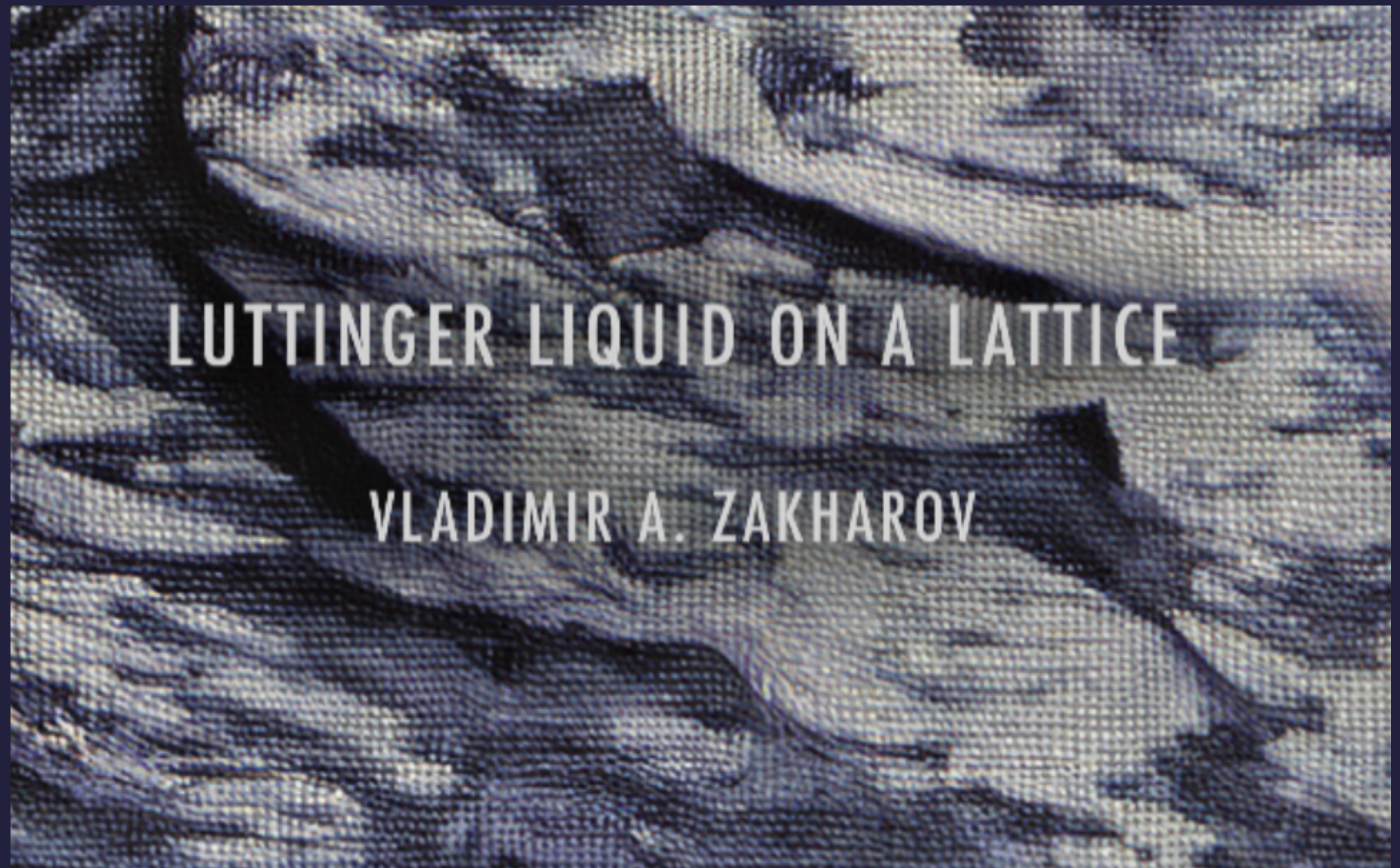
Congjun Wu,^{1,2} B. Andrei Bernevig,¹ and Shou-Cheng Zhang¹

The edge states of the recently proposed quantum spin Hall systems constitute a new symmetry class of one-dimensional liquids dubbed the “helical liquid,” where the spin orientation is determined by the direction of electron motion. We prove a no-go theorem which states that a helical liquid with an odd number of components cannot be constructed in a purely 1D lattice system. In a helical liquid with an odd number of components, a uniform gap in the ground state can appear when the time-reversal symmetry is spontaneously broken by interactions.

tangent fermions allow access to this effect on a 1D lattice

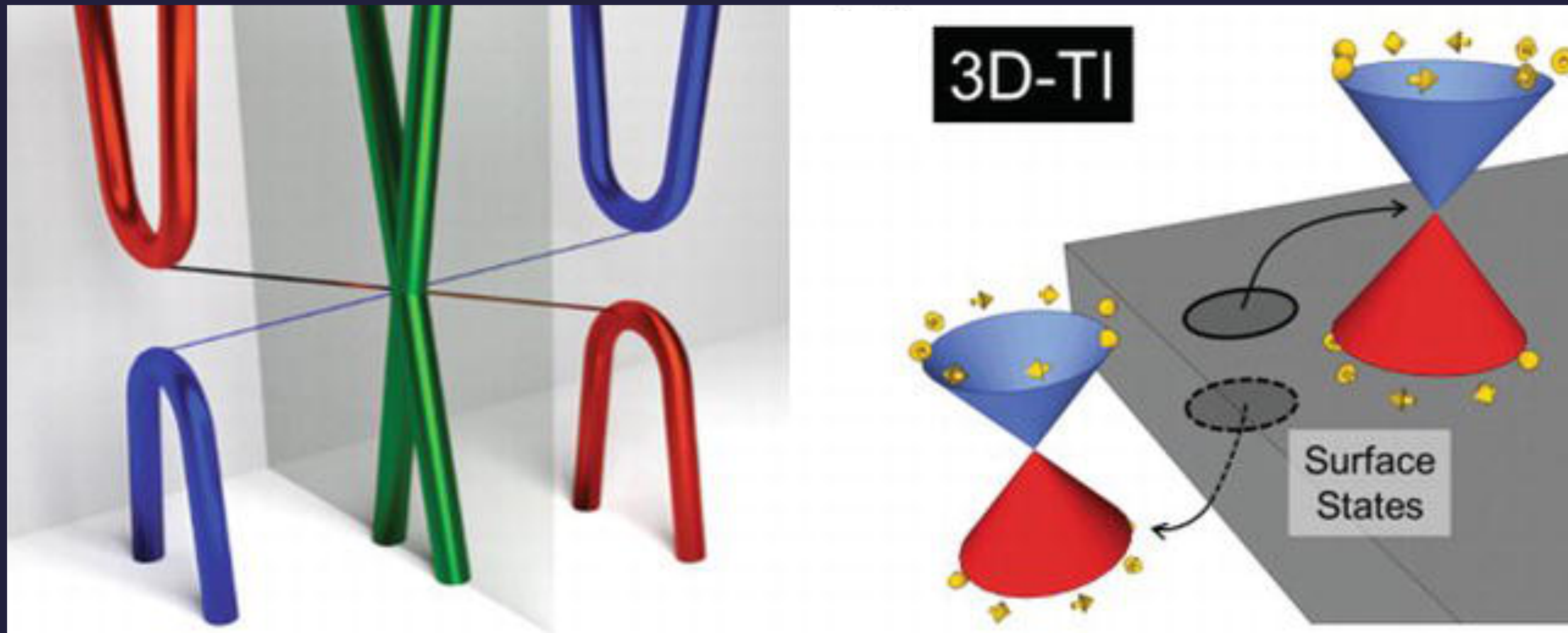
Tangent fermions:

lattice fermions with a topologically
protected Dirac cone



3D topological insulator

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insulating bulk with gapless Dirac cones on the surface, protected from localisation by time-reversal symmetry

$$H = p_x \sigma_x + p_y \sigma_y + V(x, y)$$

can we discretize H without gapping out the Dirac cone?

Fermion doubling

a classic problem in high-energy physics (lattice QCD)

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Fermion doubling

From Wikipedia, the free encyclopedia

The **fermion doubling problem** is a problem that is encountered when naively trying to put **fermionic fields** on a **lattice**. It consists in the appearance of spurious states, such that one ends up having 2^d fermionic particles (with d the number of discretized dimensions) for each original fermion. In order to solve this problem, several strategies are in use, such as **Wilson fermions** and **staggered fermions**.

Susskind fermions

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- 2 [The Nielsen–Ninomiya theorem](#)

both these methods break the symplectic TRS

How to avoid fermion doubling

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on a lattice $p \rightarrow a^{-1} \sin pa$ (periodic in $2\pi/a$), producing a Dirac cone at $p=0$ and at $p=\pi/a$
avoid the doublers by making them massive:

$$H = 2\sigma_x \tan(p/2)$$

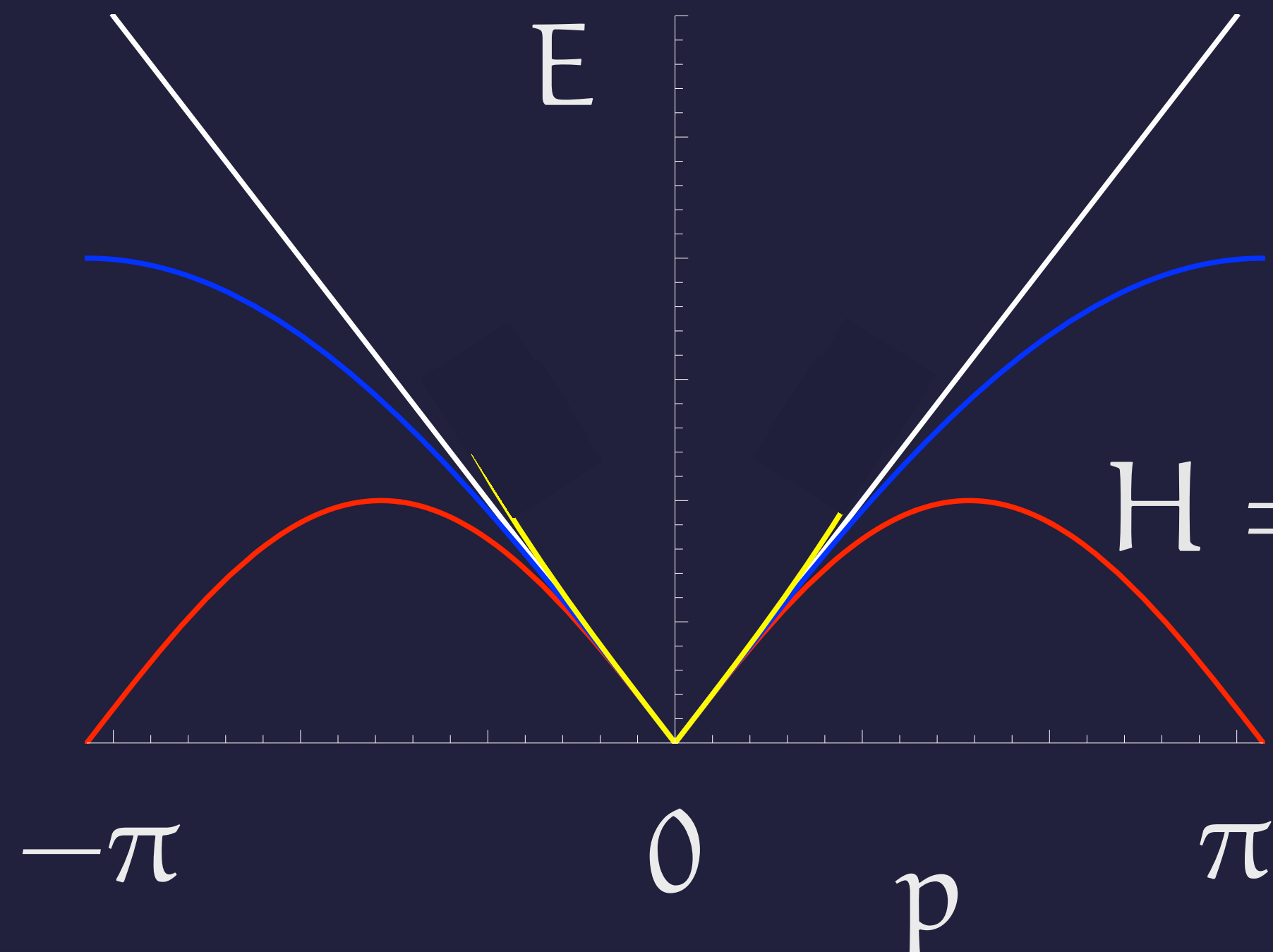
$$H = p\sigma_x$$

$$H = \sigma_x \sin p + \sigma_z(1 - \cos p)$$

Wilson or Susskind fermions – TRS broken

$$H = \sigma_x \sin p$$

tangent dispersion (Stacey, 1982) preserves TRS,
with a *nonlocal* hopping: $t_{nm} \propto (-1)^{n-m}$



Restoring locality

$$2 \tan(p/2) = \frac{\sin p}{\frac{1}{2}(1 + \cos p)}$$

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the tangent dispersion $E=2\tan(p/2)$ can be obtained from a *local* generalised eigenproblem

$$H\psi = EP\psi, \quad H = \sigma_x \sin p, \quad P = \frac{1}{2}(1 + \cos p)$$

2D generalisation: $E = 2\sqrt{\tan^2(p_x/2) + \tan^2(p_y/2)}$

$$H = \frac{1}{2}\sigma_x \sin p_x (1 + \cos p_y) + \frac{1}{2}\sigma_y \sin p_y (1 + \cos p_x)$$

$$P = \frac{1}{4}(1 + \cos p_x)(1 + \cos p_y)$$

[arXiv:2103.15615](https://arxiv.org/abs/2103.15615)

tangent fermion tensor network

Matrix-Product-Operator (MPO) representation of the Hamiltonian

$$H = \sum_{n>m=1}^N (t_{nm} c_n^\dagger c_m + \text{H.c.}) = M_1 M_2 \dots M_N$$

$t_{nm} \propto (-1)^{n-m}$ matrices M_n depend only on site n

Q: is this possible with a dimension of M_n that does not grow with N ?

arXiv:2407.06713 **YES** any dispersion of the form $P(k)\psi = Q(k)E\psi$ with P and Q polynomials in e^{ik} , has an MPO representation of fixed bond dimension

helical Luttinger liquid with Hubbard interaction

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$$H = \sum_{n>m=1}^N \left[t_{nm} (c_{n\uparrow}^\dagger c_{m\uparrow} - c_{n\downarrow}^\dagger c_{m\downarrow}) + \text{H.c.} \right] + \sum_{n=1}^N U_n$$

$$t_{nm} = 2it_0(-1)^{n-m} \quad \text{tangent dispersion:}$$

$$M_n = \begin{pmatrix} 1 & c_{n\uparrow} & c_{n\uparrow}^\dagger & c_{n\downarrow} & c_{n\downarrow}^\dagger & (2it_0)^{-1}U_n \\ 0 & -1 & 0 & 0 & 0 & c_{n\uparrow}^\dagger \\ 0 & 0 & -1 & 0 & 0 & c_{n\uparrow} \\ 0 & 0 & 0 & -1 & 0 & -c_{n\downarrow}^\dagger \\ 0 & 0 & 0 & 0 & -1 & -c_{n\downarrow} \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix},$$

exact MPO representation
with *fixed* bond dimension 6

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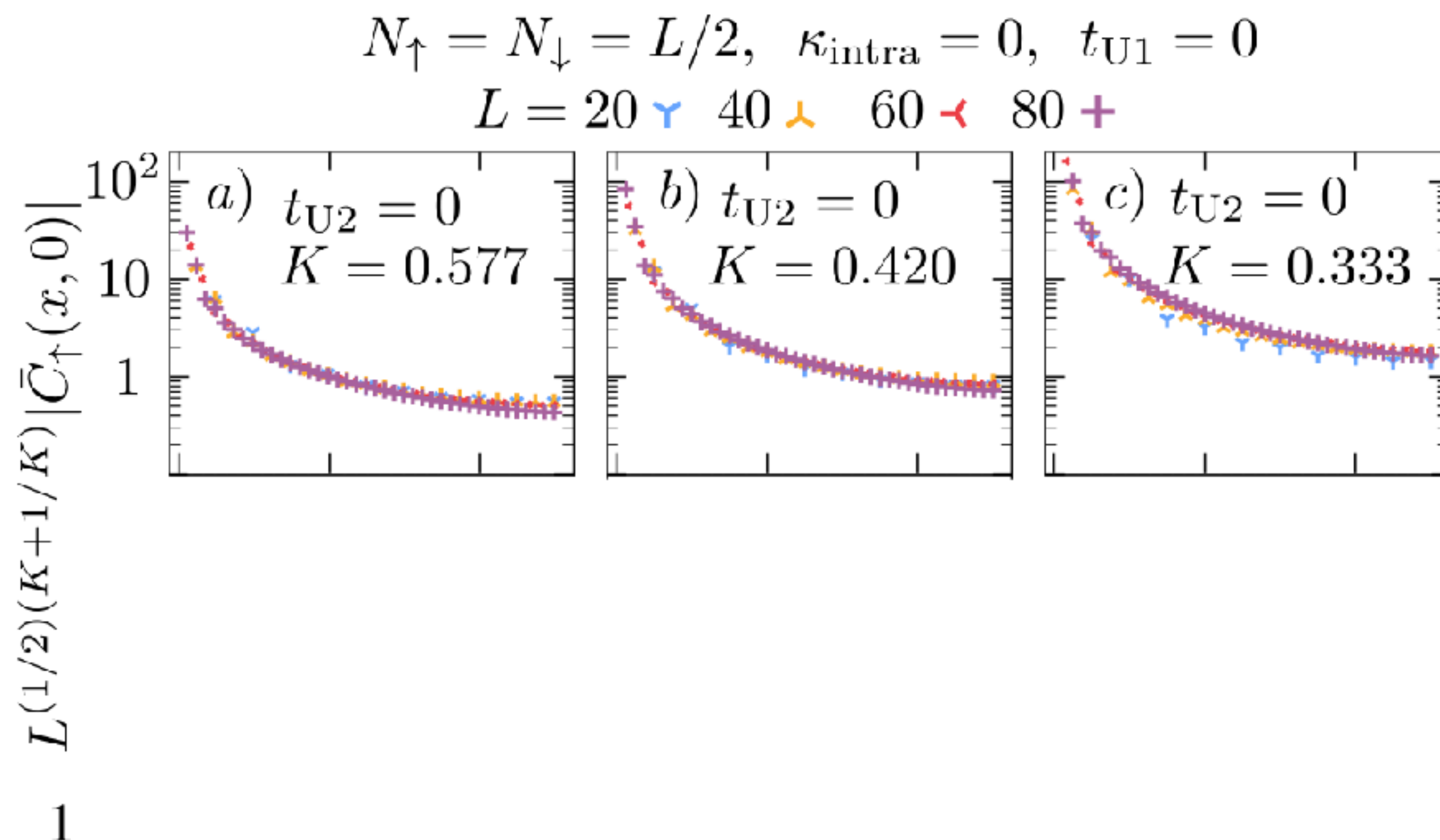
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power law scaling

$$C_{\sigma}(x, 0) = L^{-(1/2)(K+1/K)} f_C(x/L)$$

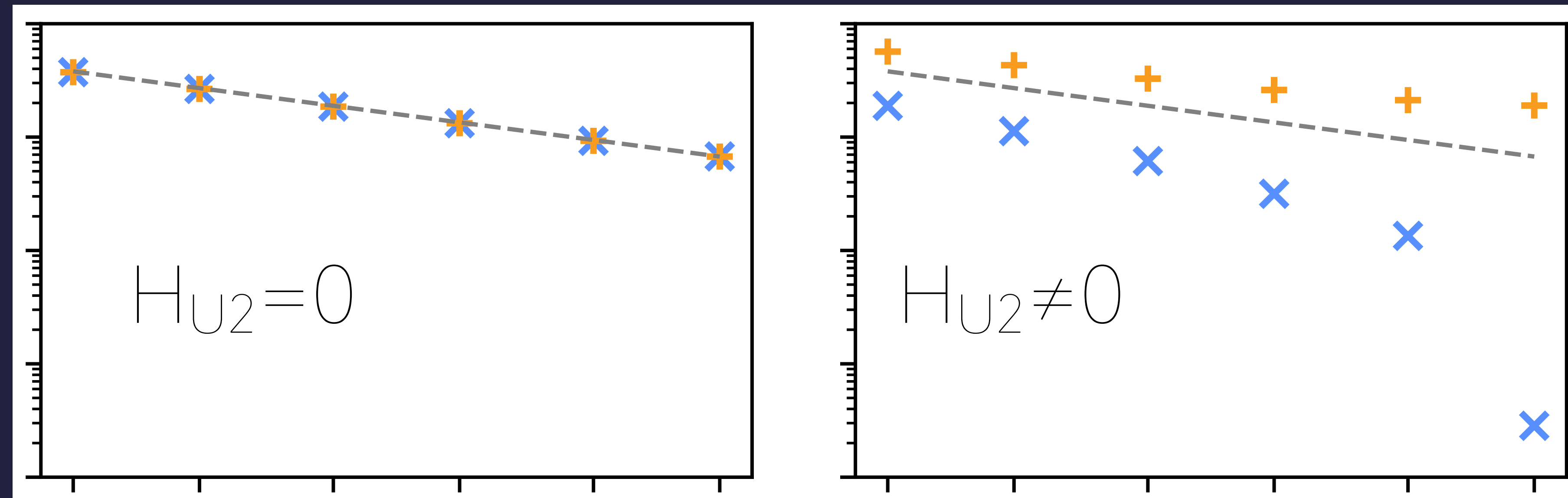
is broken by two-particle
backscattering if $K < 1/2$

$$H_{U2} = t_{U2} \sum_n [c_{n\uparrow}^{\dagger} c_{n\downarrow} c_{n+1,\uparrow}^{\dagger} c_{n+1,\downarrow} + \text{H.c.}]$$

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first two excited states as a function of system size



ground state becomes
doubly degenerate

Summary:

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- *tangent fermions*: space-time discretization of the Dirac equation without fermion doubling and without zone boundary singularities
- gapless Dirac cone protected by time-reversal & chiral symmetry
- efficient quantum Monte Carlo and tensor network implementation
- vector potential can be included in a gauge invariant way
- applications: Casimir effect, Klein tunneling, Majorana metal, **QSHE**

Tangent Fermions: Dirac or Majorana Fermions on a Lattice Without Fermion Doubling

C. W. J. Beenakker,* A. Donís Vela, G. Lemut, M. J. Pacholski, and J. Tworzydło

arXiv:2302.12793

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Lattice fermion simulation of spontaneous time-reversal symmetry breaking in a helical Luttinger liquid

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 (Dated: January 2026)

arXiv:2601.09563

The Capri Spring School on Transport in Nanostructures

Capri, April 3 - 10, 2005.

thank you so much for 20 years of memorable meetings

Lecturers

- **Carlo Beenakker**: Shot noise and full counting statistics
- **Gerrit Bauer**: Spintronics
- **Carlo Di Castro**: Low dimensional disordered conductors
- **Karsten Flensberg**: Molecular electronics
- **Andrei Komnik**: Quantum impurity problems in one-dimensional metals
- **Allan MacDonald**: Quantum Hall effect
- **Francesco Petruccione** *: Quantum information processing with nanostructures
- **Hugues Pothier**: Dephasing and relaxation in disordered nanowires
- **Arturo Tagliacozzo**: Kondo and Fano resonances
- **Amir Yacoby**: Experimental study of electrons confined to one dimension