

Coupling-energy driven pumping through quantum dots: the role of coherences

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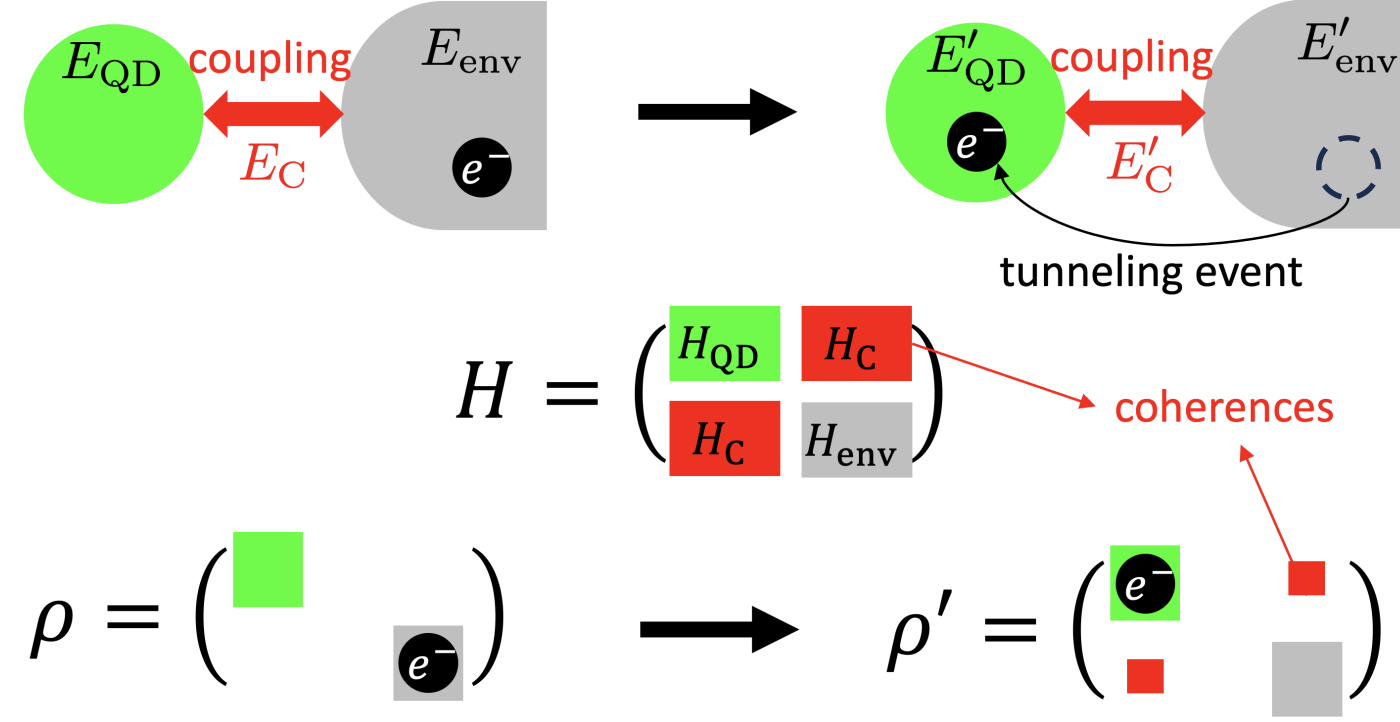


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Motivation

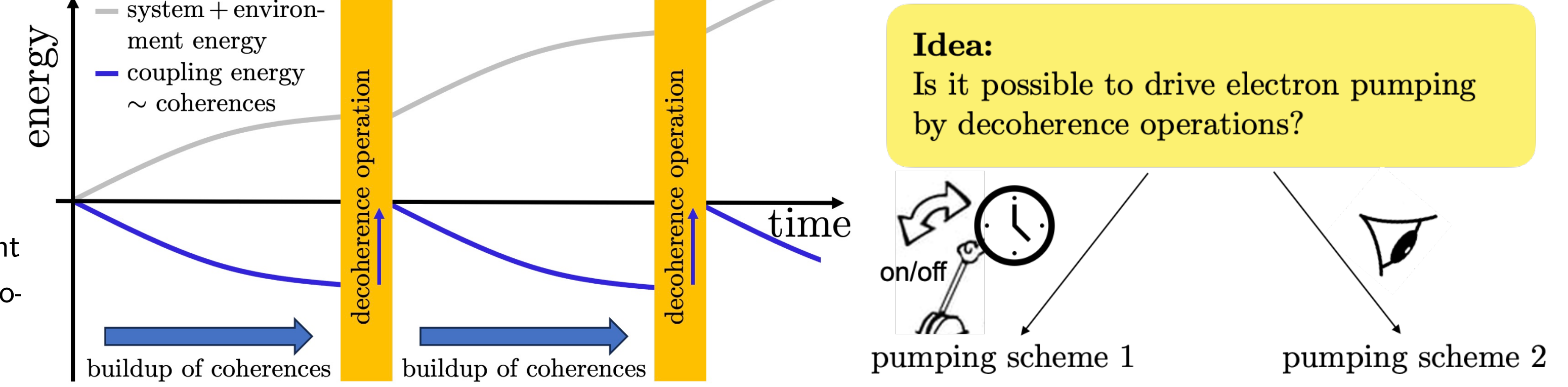
Energy conservation of a tunneling event



- coupling energy E_C has to be taken into account
 $E_{\text{QD}} + E_{\text{env}} + E_C = E'_{\text{QD}} + E'_{\text{env}} + E'_C$
- coupling energy E_C
– is related to coherences between QD and environment
– enables off-resonant and higher-order tunneling processes LL et al., PRB 111, 085103 (2025)

Destruction of quantum coherence can pump energy into a system

Schulman et al. PRL 97 240405 (2006)

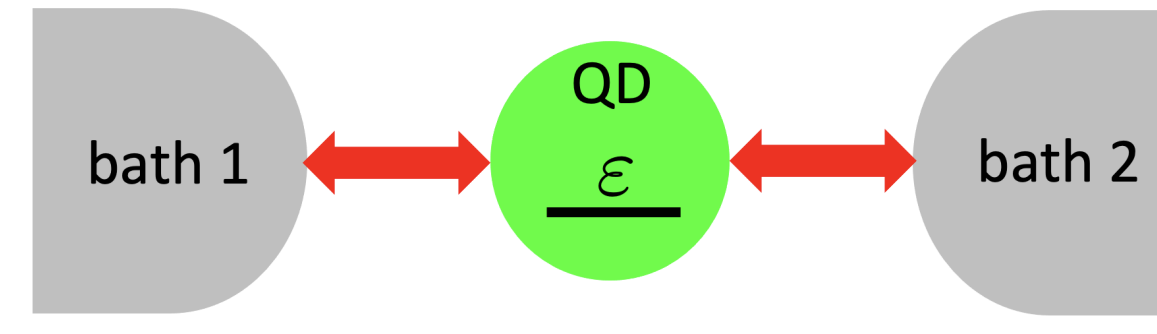


Idea:
Is it possible to drive electron pumping by decoherence operations?

Model and method

- non-interacting single-level QD coupled to two fermionic and baths:

$$H = \underbrace{\varepsilon c^\dagger c + \sum_{\nu=1}^2 \sum_k (\gamma_{\nu,k}(t) a_{\nu,k}^\dagger c + \text{h.c.})}_{\equiv H_C(t)} + \underbrace{\sum_{\nu,k} \varepsilon_{\nu,k} a_{\nu,k}^\dagger a_{\nu,k}}_{\equiv H_{\text{env}}}$$

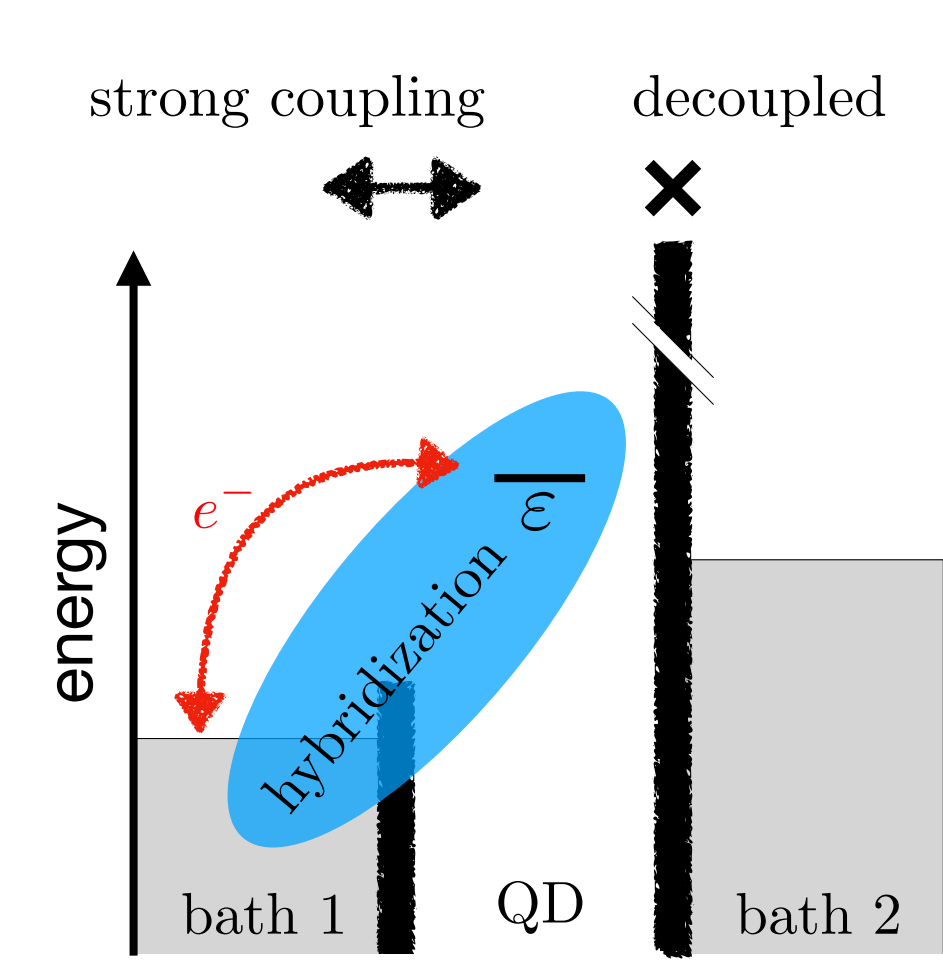


- exact solution with Heisenberg equation and Laplace transform technique Topp et al. EPL 110 67003 (2015)

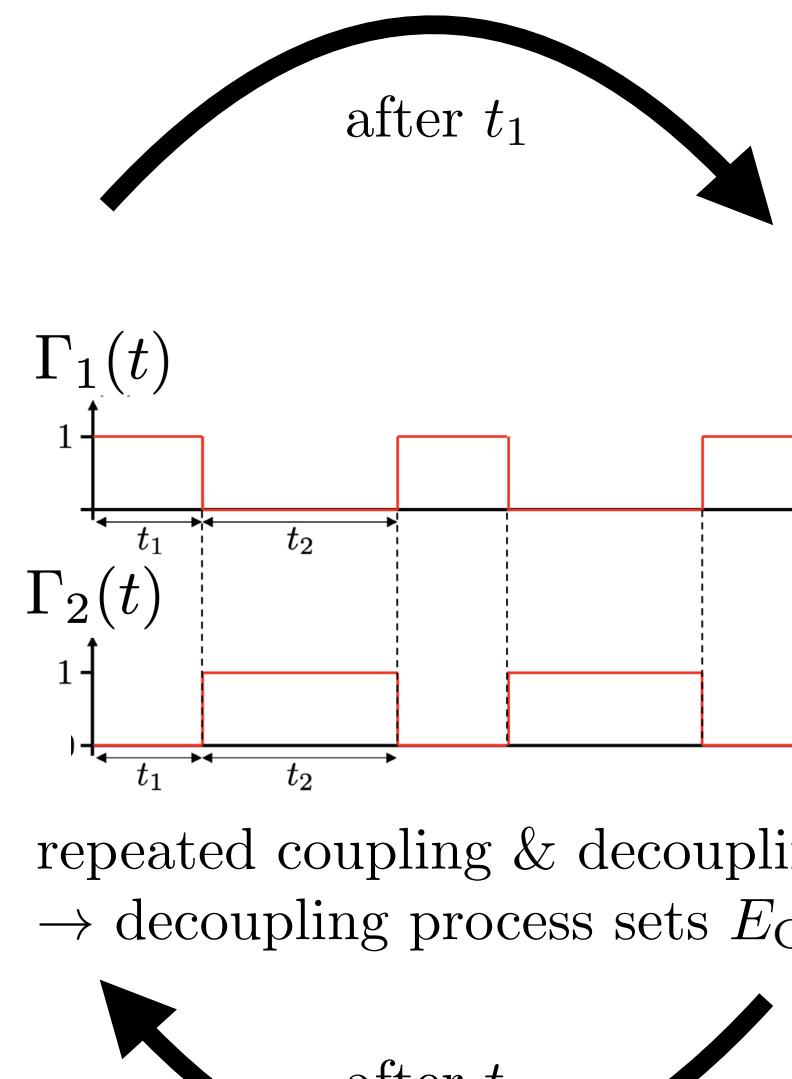
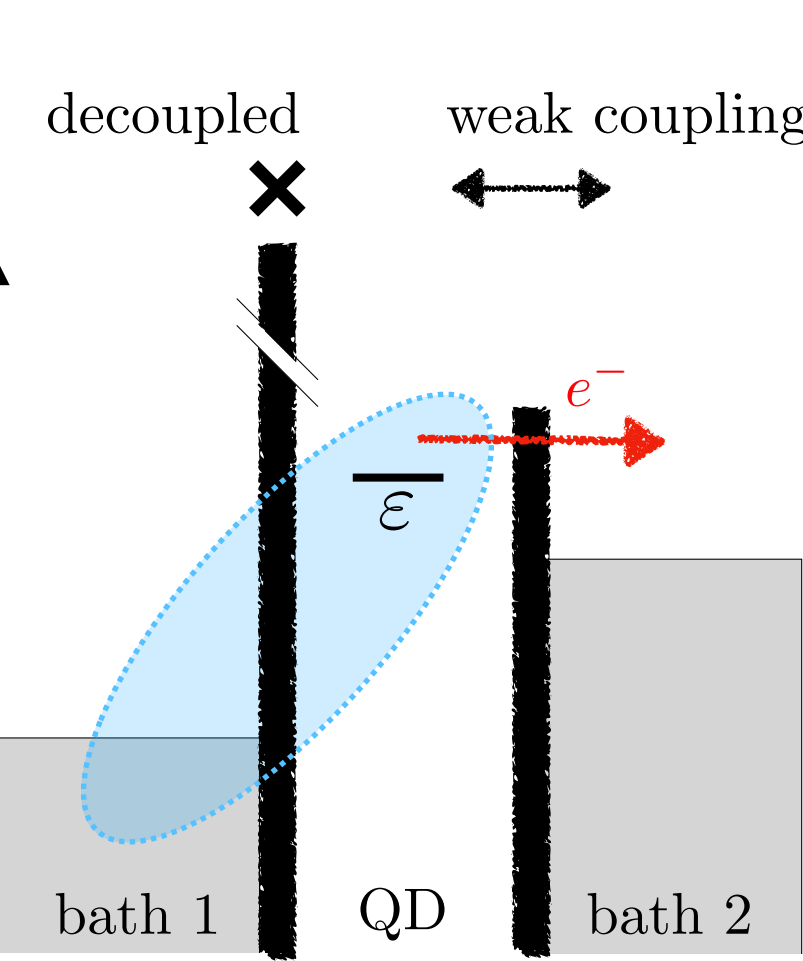
$$\left. \begin{aligned} \frac{d}{dt} c_H(t) &= \frac{i}{\hbar} [H, c_H(t)] \\ \frac{d}{dt} a_{H,\nu,k}(t) &= \frac{i}{\hbar} [H, a_{H,\nu,k}(t)] \end{aligned} \right\} \xrightarrow{\text{Laplace trans.}} \text{solve linear equations} \xrightarrow{\text{inv. Laplace trans.}} c_H(t) \text{ \& } a_{H,\nu,k}(t)$$

Pumping scheme I: Coupling/Decoupling

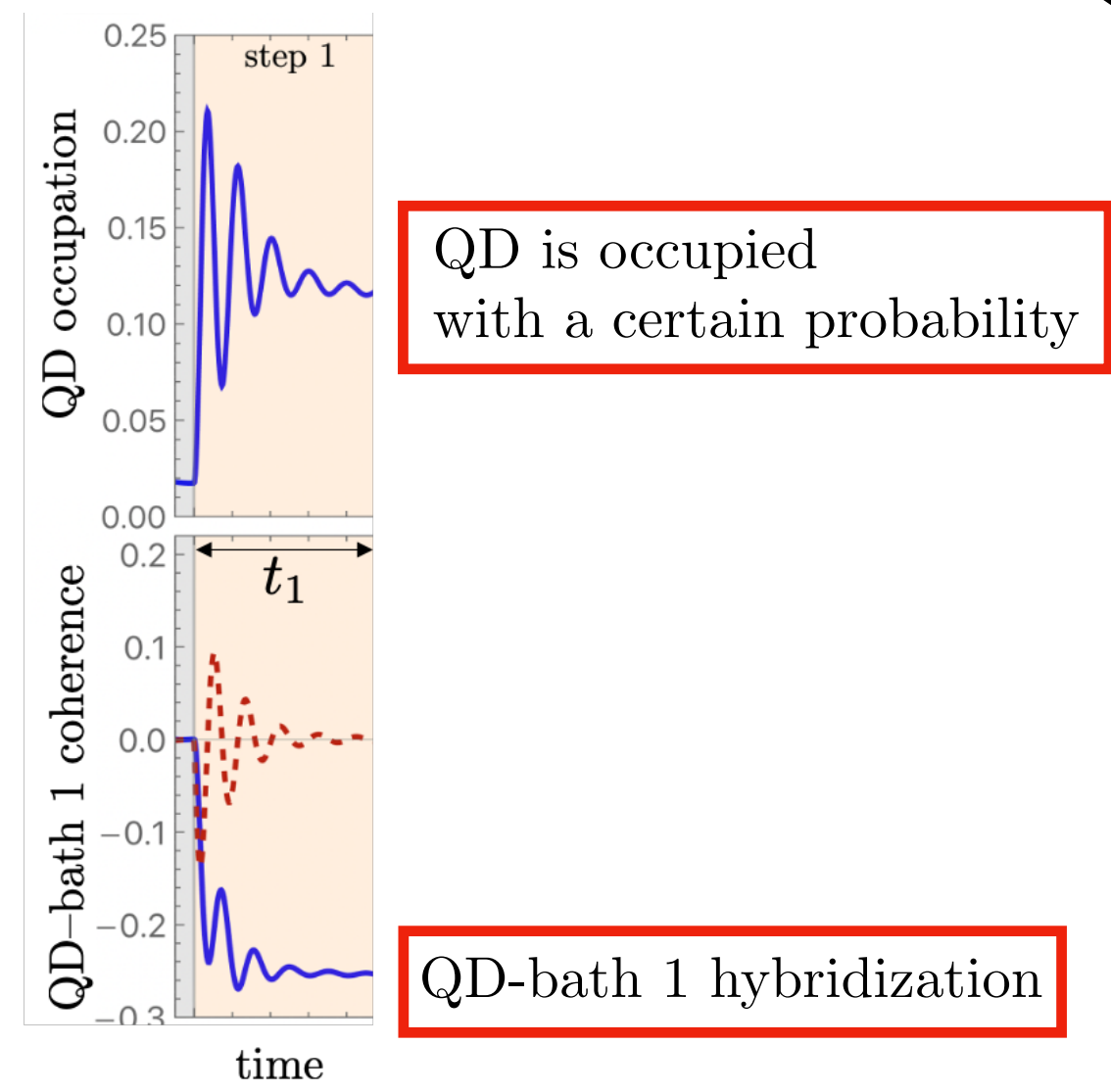
step 1: loading



step 2: unloading

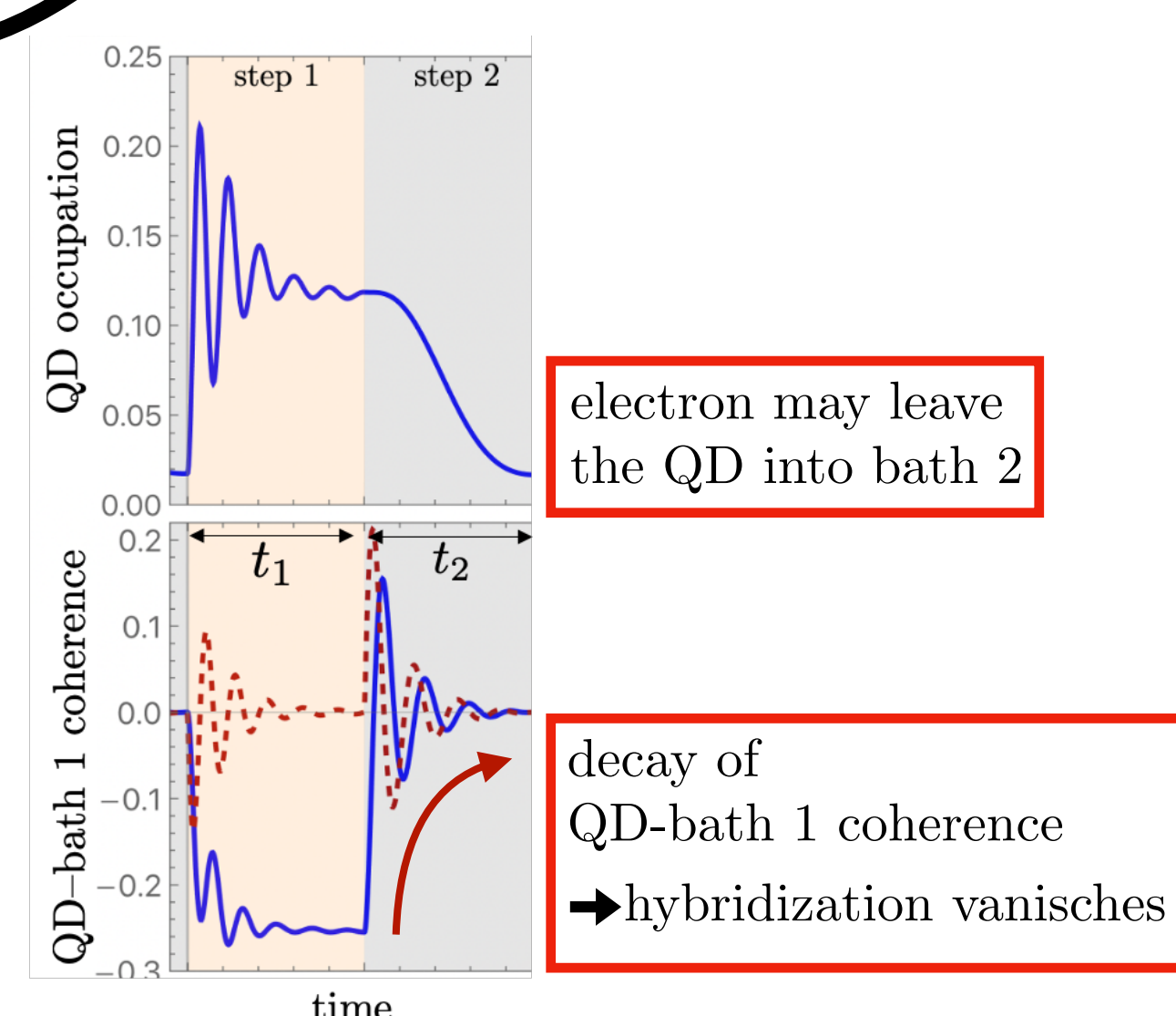


repeated coupling & decoupling
→ decoupling process sets $E_C = 0$



QD is occupied with a certain probability

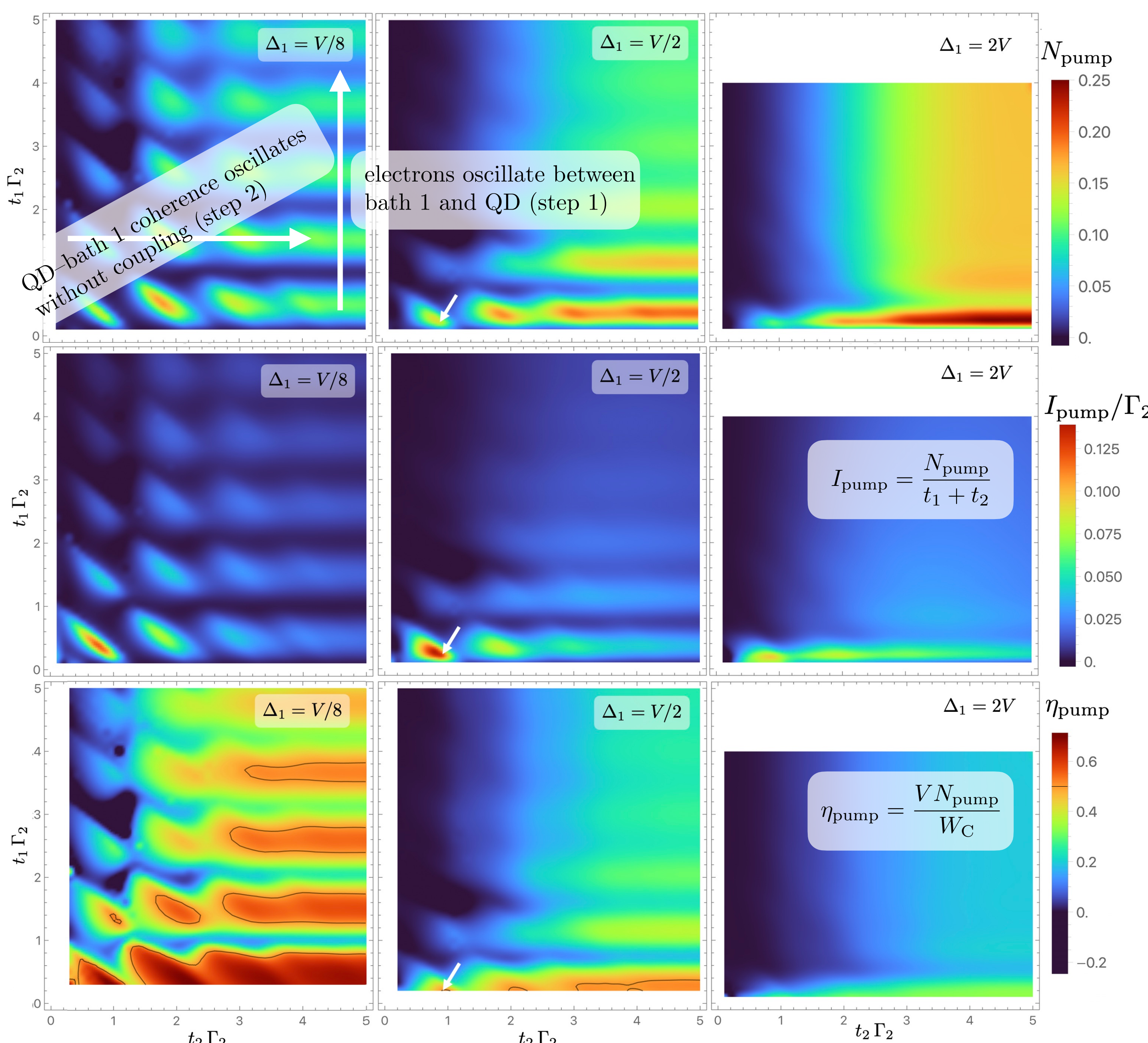
QD-bath 1 hybridization



electron may leave the QD into bath 2

decay of QD-bath 1 coherence
→ hybridization vanishes

- current I_{pump} against applied voltage if $\Gamma_1 \gg \Gamma_2 \leftrightarrow$ driven by off-resonant tunneling processes



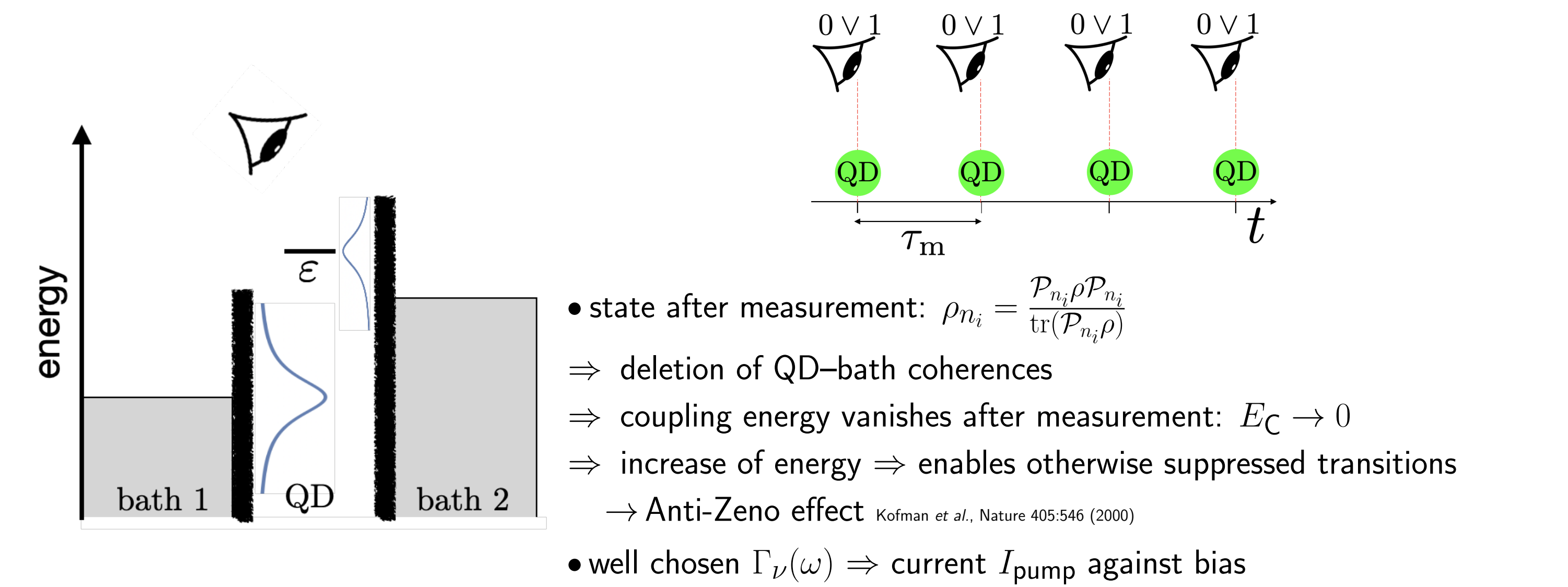
- pumping against bias

- oscillations with t_1 and t_2

Pumping scheme II: Measurement induced pumping

Projective measurements:

- stroboscopically repeated measurements of the particle number



state after measurement: $\rho_{n_i} = \frac{P_{n_i} \rho P_{n_i}}{\text{tr}(P_{n_i} \rho)}$

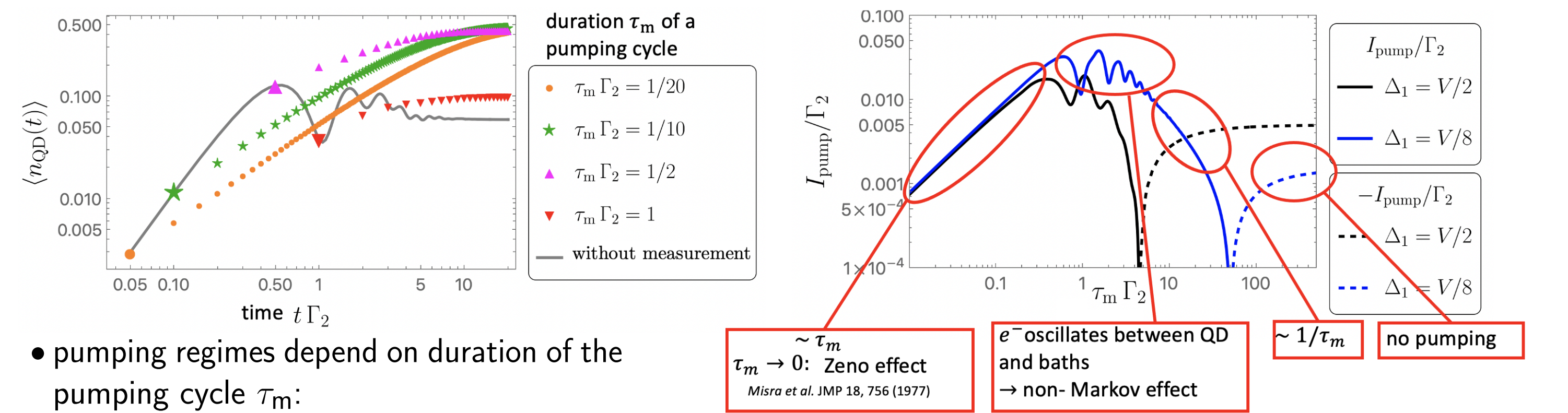
⇒ deletion of QD-bath coherences

⇒ coupling energy vanishes after measurement: $E_C \rightarrow 0$

⇒ increase of energy ⇒ enables otherwise suppressed transitions

→ Anti-Zeno effect Kofman et al., Nature 405:546 (2000)

- well chosen $\Gamma_\nu(\omega) \Rightarrow$ current I_{pump} against bias



- pumping regimes depend on duration of the pumping cycle τ_m :

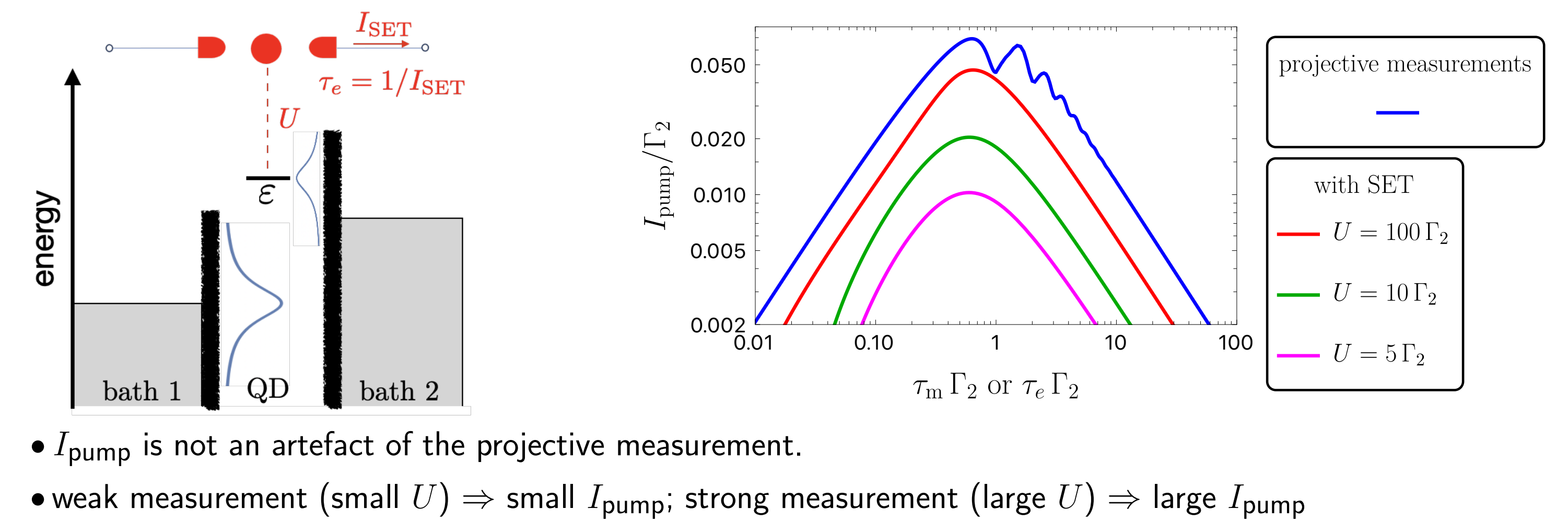
$\tau_m \rightarrow 0$: Zeno effect
Mitra et al. JMP 18, 756 (1977)

e^- oscillates between QD and baths
→ non-Markov effect

$\sim 1/\tau_m$ no pumping

Weak measurement via a single-electron transistor (SET):

- projective measurement is unrealistic → simulating weak continuous measurement via SET
- Redfield master-equation + fermionic reaction-coordinates mapping

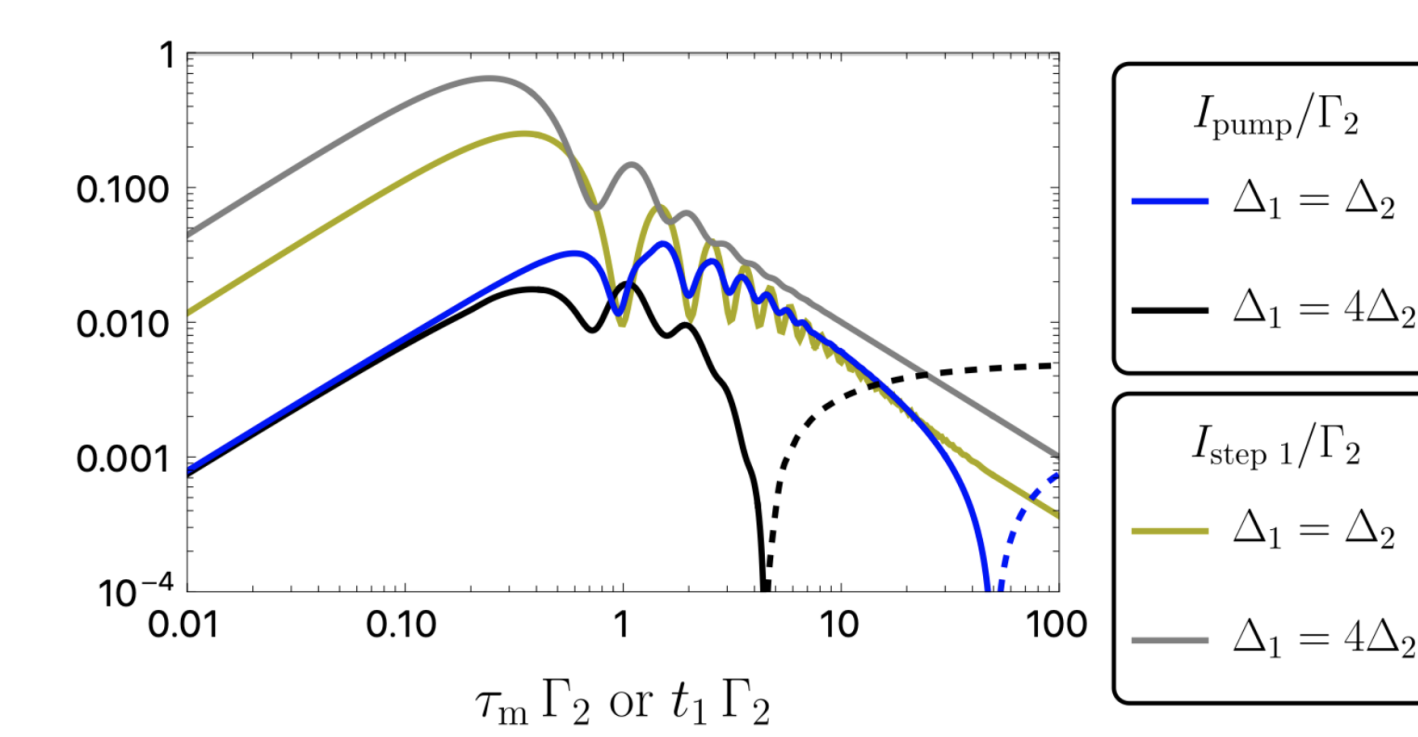


- I_{pump} is not an artefact of the projective measurement.

- weak measurement (small U) ⇒ small I_{pump} ; strong measurement (large U) ⇒ large I_{pump}

Conclusion

- pumping against bias is possible
– both schemes rely on the same mechanisms → “pure quantum-effect”
- Non-Markovian effects enhance efficiency
– clear signatures of quantum coherences



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