

Higher order boundary signatures of crystalline symmetry protected topological phases

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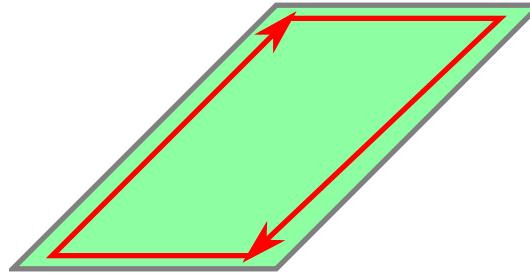
Capri, 09.05.2019

Topological phases of fermions

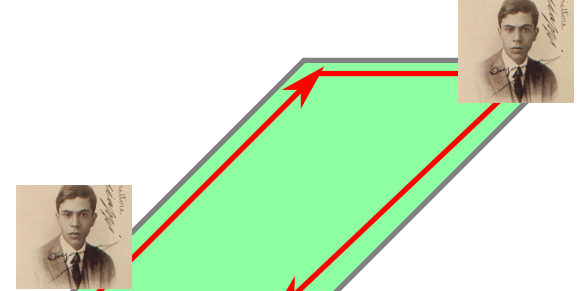
... that can be described by a quadratic (mean field) theory



1d topological superconductor:
Kitaev Wire

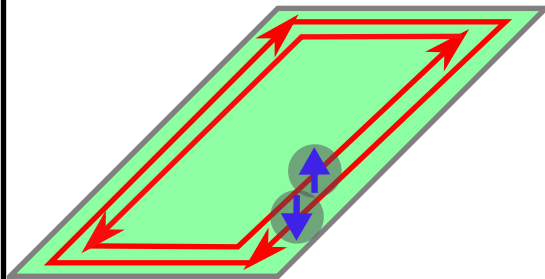


Quantum Hall effect

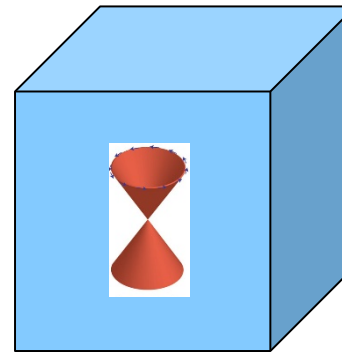


$p_x + ip_y$ wave superconductor

Topological phases protected by a local symmetry: Time reversal



Quantum
spin-Hall effect

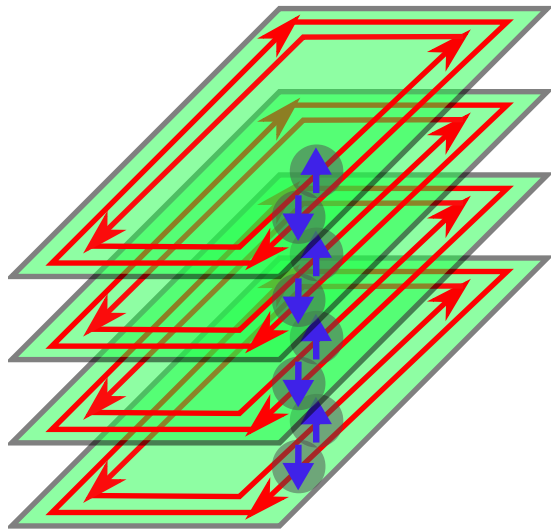


3d topological
insulator (TI)

Bulk boundary correspondence
gapped topological bulk = gapless **anomalous** states on a symmetric boundary

Crystalline symmetry protected topological phases

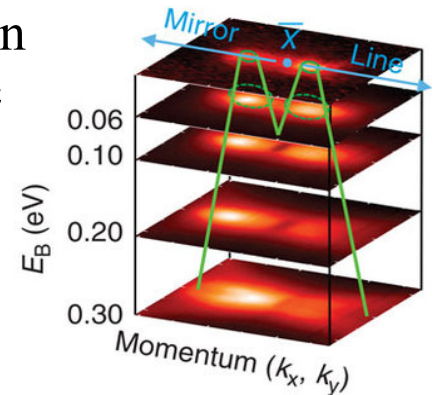
Translation symmetry:
Weak topological insulators



Only boundaries that preserve the protecting translation symmetry are gapless

Mirror symmetry: SnTe

Double Dirac cone on reflection-symmetric surface

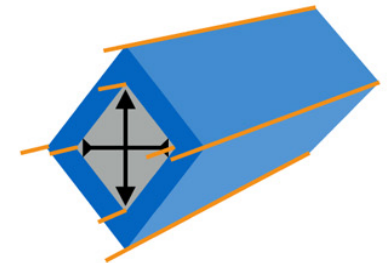


Hsieh *et al.* (2012)

Xu *et al.* (2012)

Tanaka *et al.* (2012)

Helical hinge modes in a strained lattice

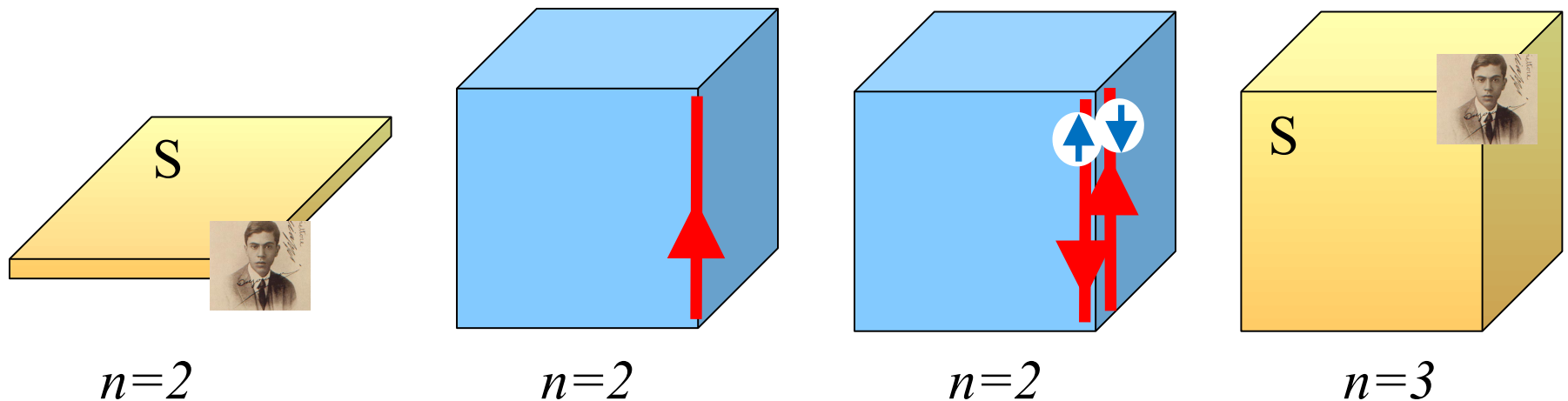


Schindler *et al.* (2018)

Crystalline topological phases with higher order boundary signatures

d-dimensional n th order crystalline SPT:

- Anomalous gapless states form a $(d-n)$ dimensional manifold on the boundary



Signatures of “purely crystalline symmetry protected topological phases”

→ Trivial when crystalline symmetry is broken

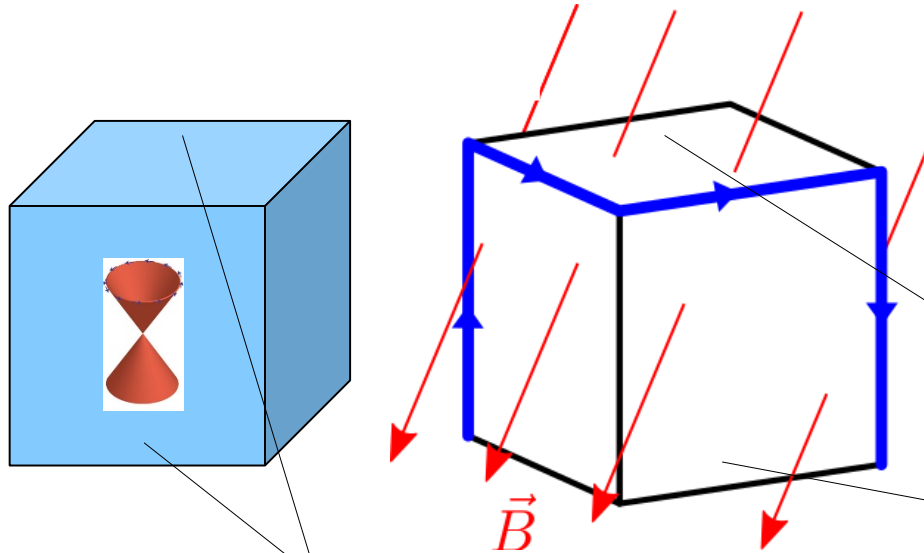
Not every higher order boundary signature corresponds to a crystalline SPT.

Which higher order boundary signatures correspond to a higher order crystalline SPT?

Extrinsic higher order boundary signatures

Example:

Gapless hinge state of a 3d TI in a magnetic field...



Surface Dirac Hamiltonian

$$H_{\text{sf}}(\vec{k}) = k_x \sigma_1 + k_y \sigma_2$$

Crossing protected by TRS: $\mathcal{T} = i\sigma_2 K$

Surface Dirac Hamiltonian

$$H_{\text{sf}}(\vec{k}) = k_x \sigma_1 + k_y \sigma_2 + \vec{B} \cdot \vec{\sigma}$$

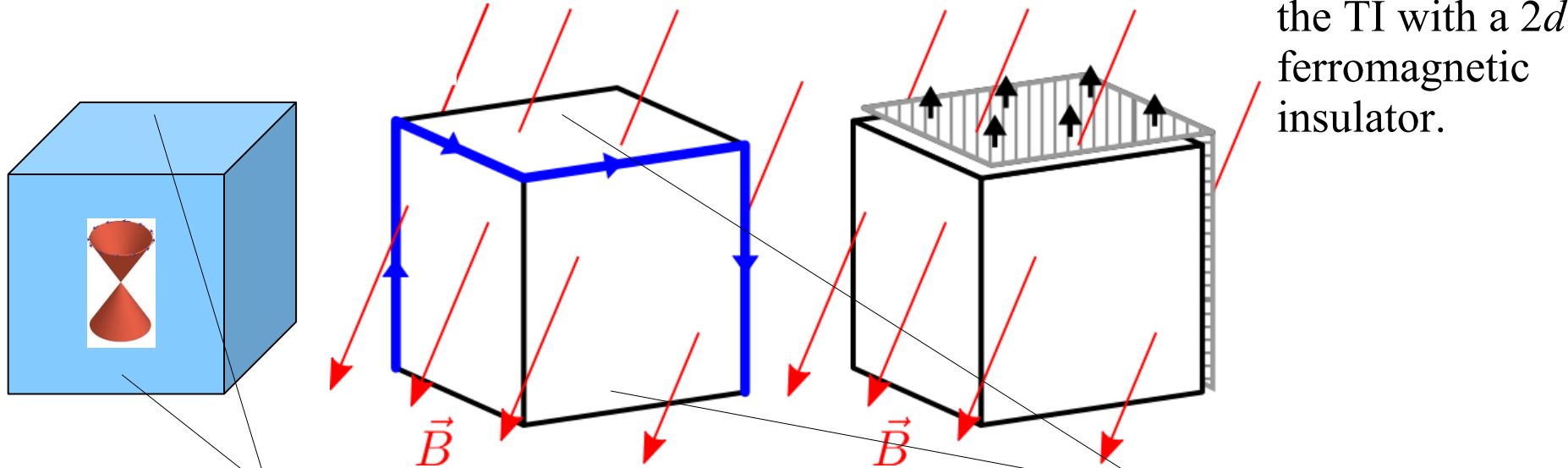
The component $B_z \sigma_z$ normal to the surface gaps the Dirac Hamiltonian!

$$\sigma_{xy} = \pm \frac{1}{2} \frac{e^2}{h}$$

Extrinsic higher order boundary signatures

Example:

Gapless hinge state of a 3d TI in a magnetic field... ... that can be removed by decorating the TI with a 2d ferromagnetic insulator.



Surface Dirac Hamiltonian

$$H_{\text{sf}}(\vec{k}) = k_x \sigma_1 + k_y \sigma_2$$

Crossing protected by TRS: $\mathcal{T} = i\sigma_2 K$

Surface Dirac Hamiltonian

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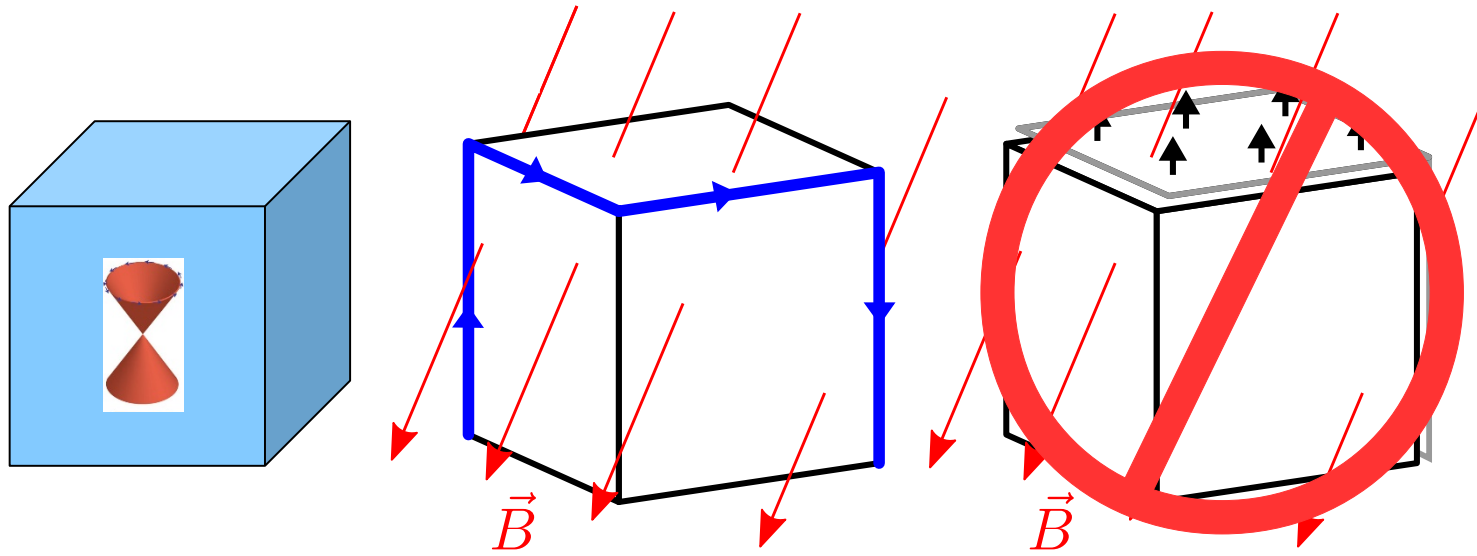
The component $B_z \sigma_z$ normal to the surface gaps the Dirac Hamiltonian!

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Extrinsic boundary signature: Existence depends on termination.

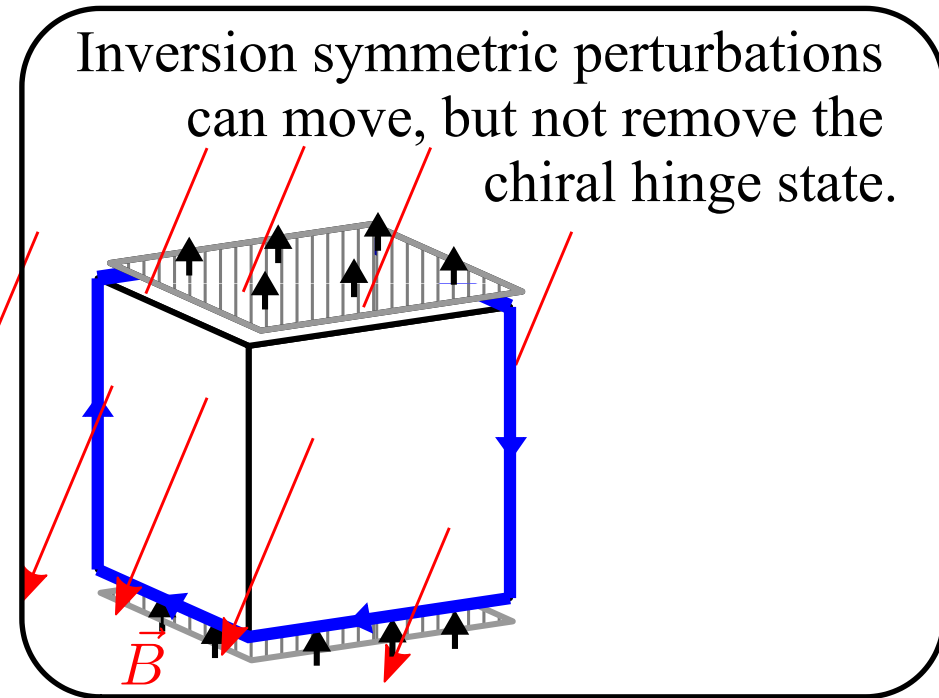
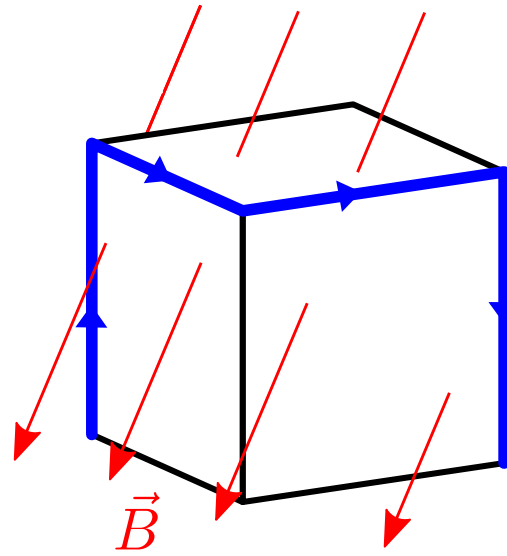
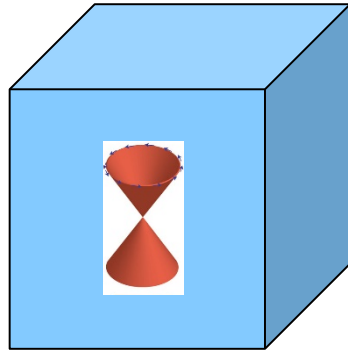
Intrinsic higher order boundary signatures

Require **inversion symmetry both in the crystal and for the sample!**



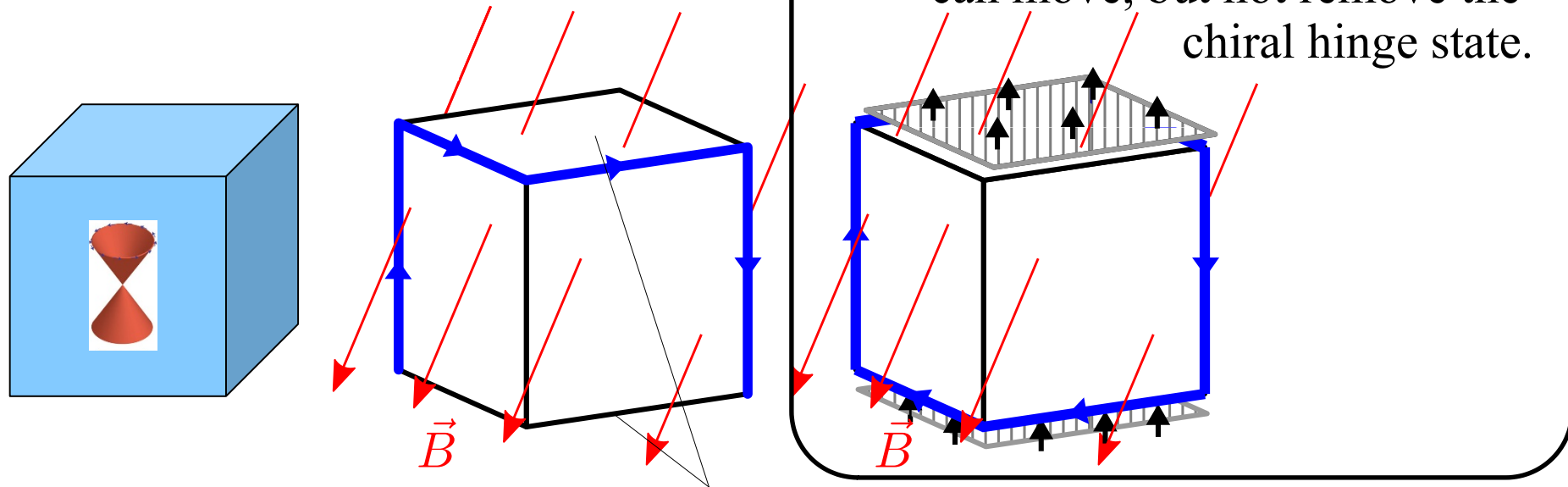
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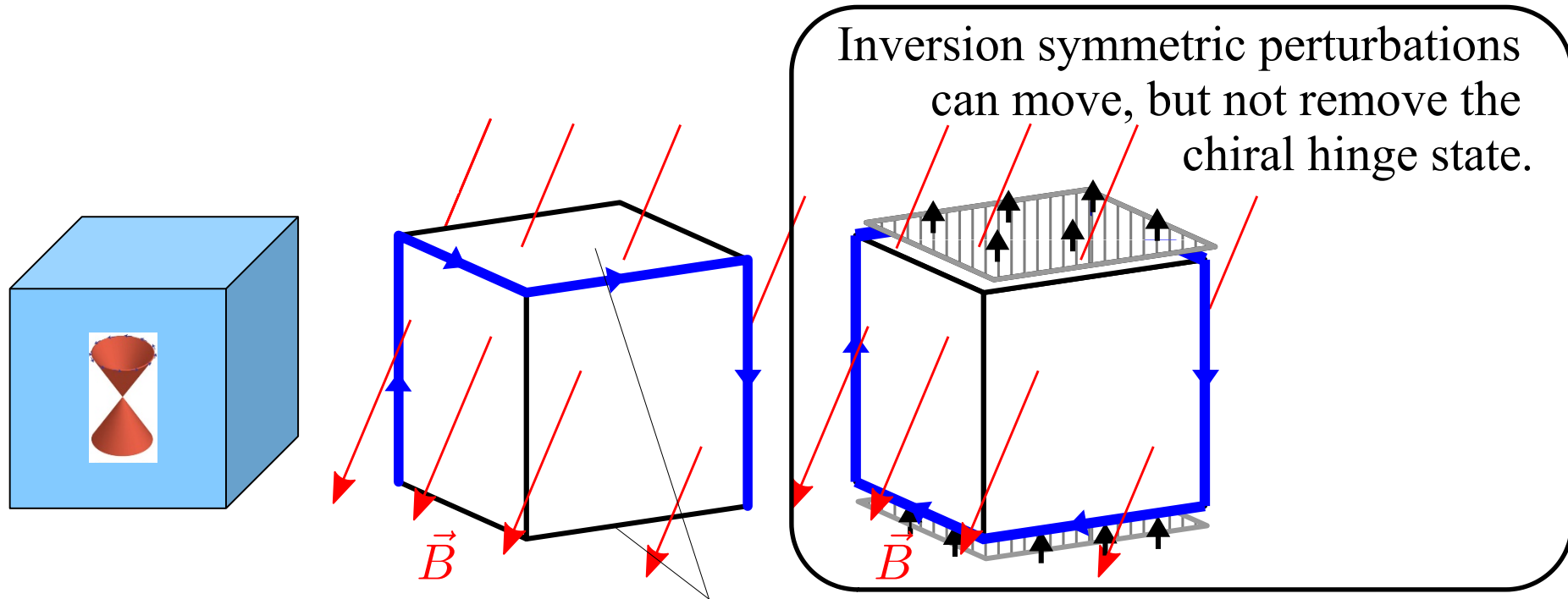
Inversion symmetry requires that inversion symmetry related surfaces have opposite mass!

$$H_{\text{sf}}(\vec{k}) = k_x \sigma_1 + k_y \sigma_2 + \vec{B} \cdot \vec{\sigma}$$

Intrinsic boundary signature: Existence independent of termination

Intrinsic higher order boundary signatures

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Intrinsic boundary signature: Existence independent of termination

Higher order bulk boundary correspondence:

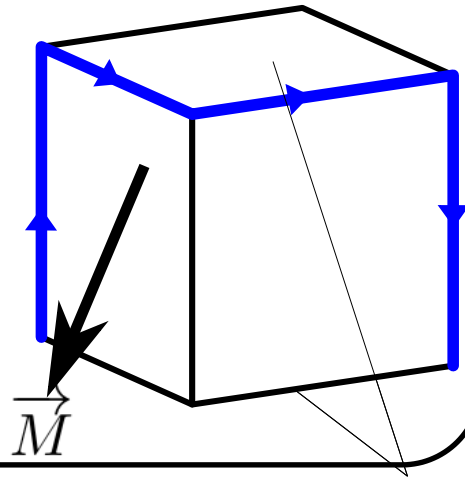
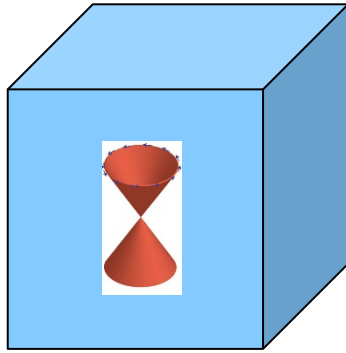
Intrinsic boundary signatures can not be induced / removed by adding any lower dimensional SPT \rightarrow Boundary signature of a topological bulk.

Intrinsic higher order boundary signatures

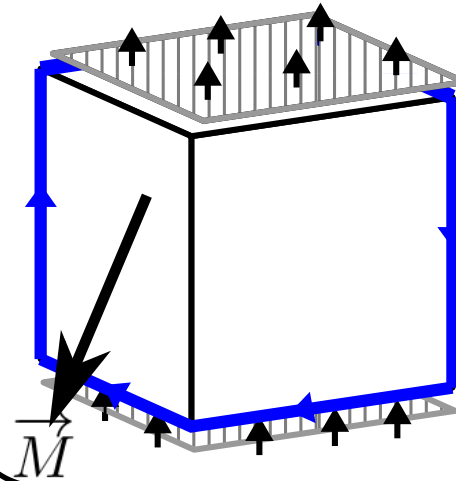
Require **inversion symmetry** both in the crystal and for the sample!

“Magnetic topological insulator”

Cr, V doped $(\text{Bi, Sb})_2\text{Te}_3$



Inversion symmetric perturbations
can move, but not remove the
chiral hinge state.



Inversion symmetry requires that inversion symmetry related surfaces have opposite mass!

$$H_{\text{sf}}(\vec{k}) = k_x \sigma_1 + k_y \sigma_2 + \vec{M} \cdot \vec{\sigma}$$

Intrinsic boundary signature: Existence independent of termination

Higher order bulk boundary correspondence:

Intrinsic boundary signatures can not be induced / removed by adding any
lower dimensional SPT \rightarrow Boundary signature of a topological bulk.

Which higher order boundary signatures correspond to a higher order crystalline SPT?

Intrinsic higher order boundary signatures:

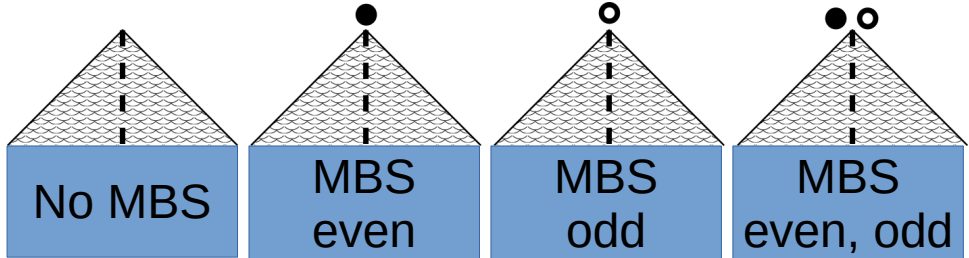
Anomalous gapless states on the surface of a crystal

- whose shape respects the space group of the crystal (up to translations)
- that can not be removed by any change of symmetric change of termination

Second order topological superconductor with mirror symmetry: Boundary perspective

Extrinsic classification of a mirror symmetric corner

$$\mathcal{K}_e = \left\{ \begin{array}{c} \text{No MBS} \\ \text{MBS even} \\ \text{MBS odd} \\ \text{MBS even, odd} \end{array} \right\}$$



p-wave / spinless SC:
 In Majorana basis:
 $H(k_1, k_2) = -H^*(-k_1, -k_2)$
 $= \sigma_1 H(-k_1, k_2) \sigma_1$

Second order topological superconductor with mirror symmetry: Boundary perspective

Extrinsic classification of a mirror symmetric corner

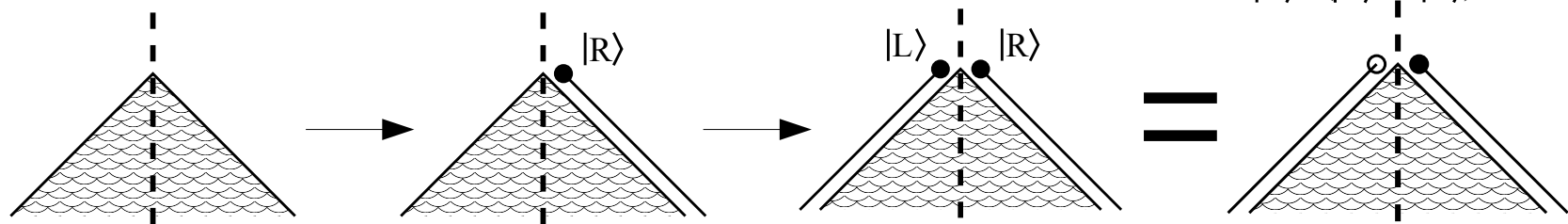
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Intrinsic classification of a mirror symmetric corner

Construct a mirror symmetric boundary decoration:



Second order topological superconductor with mirror symmetry: Boundary perspective

Extrinsic classification of a mirror symmetric corner

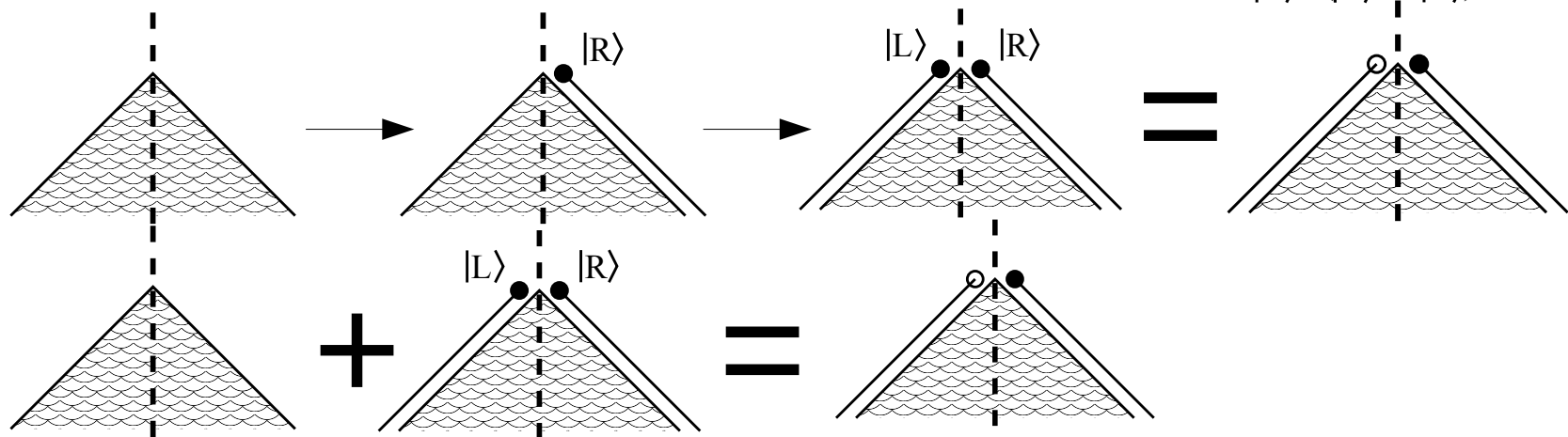
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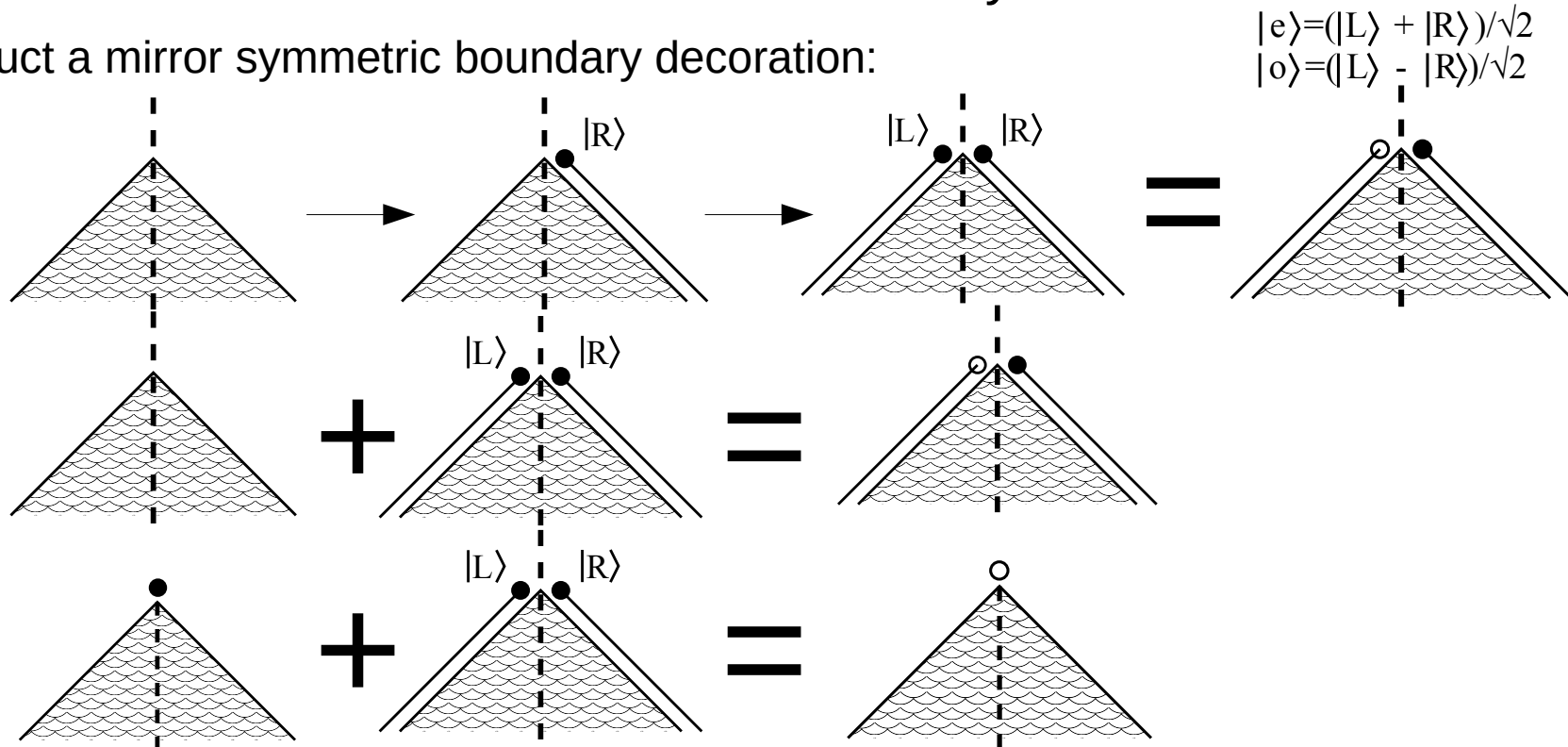
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Intrinsic classification of a mirror symmetric corner

$$\mathcal{K}_i = \mathcal{K}_e \text{ / } \begin{array}{c} \bullet \circ \\ \text{---} \end{array} = \left\{ \begin{array}{c} \text{---} \\ \bullet \text{ or } \circ \end{array} \right\}$$

Second order topological superconductor with mirror symmetry: Boundary perspective

Extrinsic classification of a mirror symmetric corner

$$\mathcal{K}_e = \left\{ \begin{array}{c} \text{No MBS} \\ \text{MBS even} \\ \text{MBS odd} \\ \text{MBS even, odd} \end{array} \right\}$$

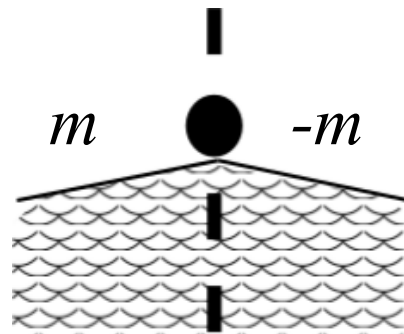
p-wave / spinless SC:
In Majorana basis:

$$H(k_1, k_2) = -H^*(-k_1, -k_2) = \sigma_1 H(-k_1, k_2) \sigma_1$$

Intrinsic classification of a mirror symmetric corner

$$\mathcal{K}_i = \mathcal{K}_e \text{ (with a diagonal line) } = \left\{ \begin{array}{c} \text{or} \end{array} \right\}$$

Bulk-boundary-corner correspondence



$$i\partial_x \sigma_3 + m(x) \sigma_2$$

$$m(x) = -m(-x)$$

Second-order topological insulators and superconductors with an order-two crystalline symmetry

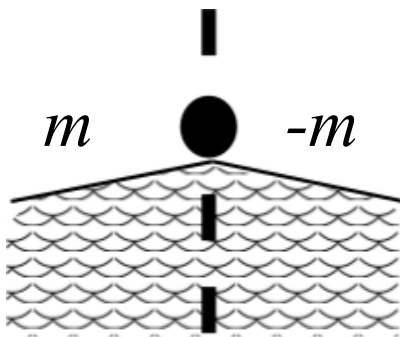
Max Geier,^{1,*} Luka Trifunovic,¹ Max Hoskam,^{1,2} and Piet W. Brouwer¹

¹*Dahlem Center for Complex Quantum Systems and Physics Department, Freie Universität Berlin, Arnimallee 14, 14195 Berlin, Germany*

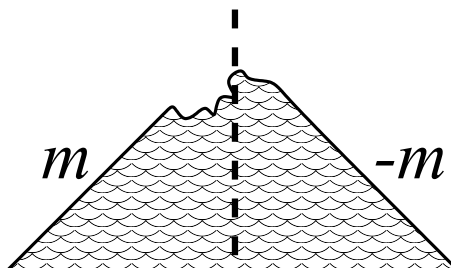
²*Department of Applied Physics, Eindhoven University of Technology, 5600 MB Eindhoven, The Netherlands*

Mirror symmetry

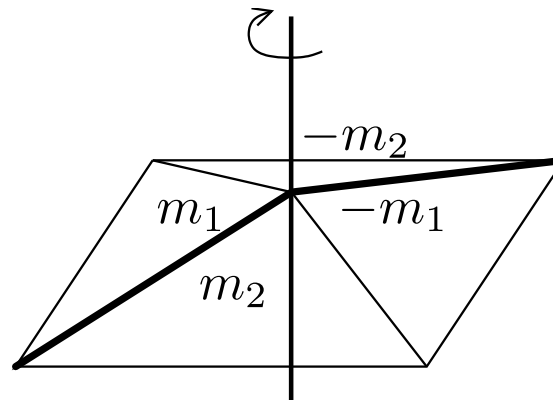
Second order
boundary signatures
without



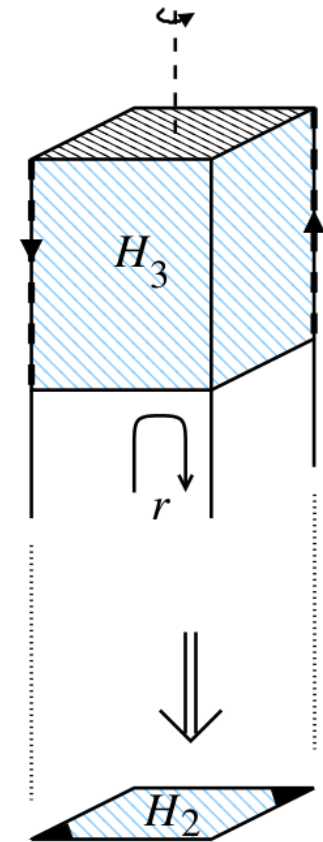
and with
symmetry breaking
at a corner



Twofold rotation symmetry



Inversion symmetry



Reflection matrix
dimensional reduction

What happens when including interactions?

Within perturbation theory:

Can the gapless boundary modes be gapped by interactions?

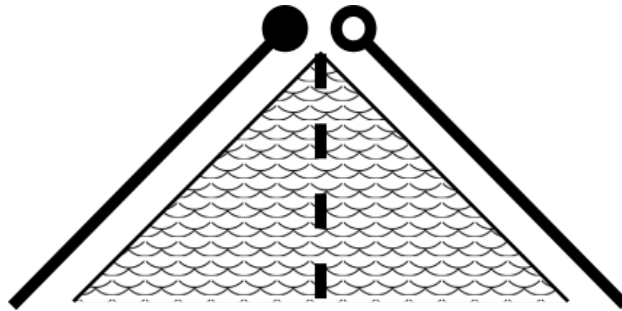
\mathbb{Z}_2 : Majorana, helical mode, ... \leftarrow NO

\mathbb{Z} : Majorana in a spinless TRS superconductor, chiral edge mode, ...

Reduce \mathbb{Z} to \mathbb{Z}_8 !

Beyond perturbation theory:

+ new phases that can not be described with a quadratic theory



Thank you for your attention :)

