

Non-local transport of strongly-interacting fermions via topological and Shockley edge states

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April 12, 2016

Scientific Reports 6, 22562 (2016)

► Floquet engineered systems

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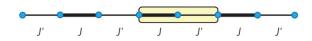
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- Haldane model: G. Jotzu et al., Nature **515** 237-240 (2014)
- Harper-Hofstadter model: M. Aidelsburguer et al., Nat. Phys.
 11 162-166 (2015)

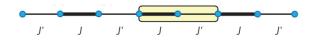
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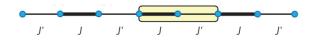


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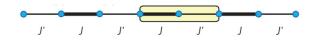
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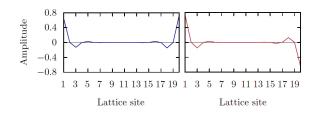
If $\lambda > 1 \Rightarrow$ Trivial phase

If $\lambda < 1 \Rightarrow$ Topological phase (Edge states)

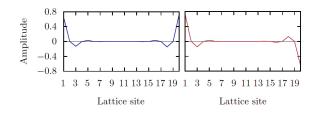
P. Delplace, PRB 84, 195452

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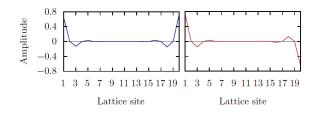


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Edge states form a non-local two-level system

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Edge states form a non-local two-level system $\Rightarrow \frac{\text{Long-range}}{\text{transfer}}$

$$H = -J' \sum_{i=1}^{M} c_{2i\sigma}^{\dagger} c_{2i-1\sigma} - J \sum_{i=1}^{M-1} c_{2i+1\sigma}^{\dagger} c_{2i\sigma} + H.c. + U \sum_{i=1}^{2M} n_{i\uparrow} n_{i\downarrow}$$

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On-site interaction

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▶ Strongly interacting regime: U > 4J

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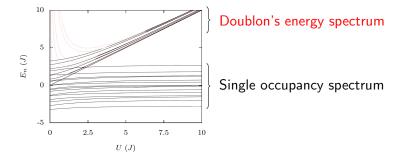
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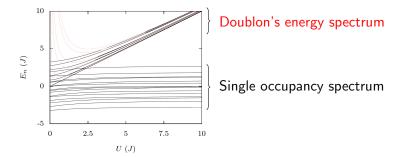
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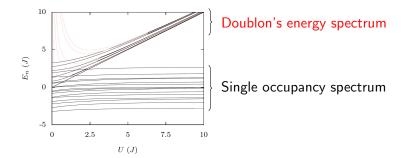
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- C. E. Creffield & G. Platero, PRB 69, 165312; PRL 105, 086804
- K. Winkler et al., Nature 441, 853-856
- S. Wall et al., Nat. Phys. 7, 114-118





We want to project out states we are not interested in



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- ▶ Unitary transformation that block-diagonalizes the Hamiltonian, perturbatively in powers of J/U and J'/U up to first order.

$$H_{\text{eff}} = J'_{\text{eff}} \sum_{i=1}^{M} d_{2i}^{\dagger} d_{2i-1} + J_{\text{eff}} \sum_{i=1}^{M-1} d_{2i+1}^{\dagger} d_{2i} + H.c. + \sum_{i=1}^{2M} \mu_i n_i^d$$

$$H_{ ext{eff}} = J_{ ext{eff}}' \sum_{i=1}^{M} d_{2i}^{\dagger} d_{2i-1} + J_{ ext{eff}} \sum_{i=1}^{M-1} d_{2i+1}^{\dagger} d_{2i} + H.c. + \sum_{i=1}^{2M} \mu_{i} n_{i}^{d}$$
 dimer tight-binding

$$\begin{aligned} H_{\text{eff}} &= J_{\text{eff}}' \sum_{i=1}^{M} d_{2i}^{\dagger} d_{2i-1} + J_{\text{eff}} \sum_{i=1}^{M-1} d_{2i+1}^{\dagger} d_{2i} + H.c. + \sum_{i=1}^{2M} \mu_{i} n_{i}^{d} \\ & \text{dimer tight-binding} \end{aligned}$$
 effective chemical potential

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Effective hoppings:
$$J'_{\rm eff}=2J'^2/U$$
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$$\mu_i = U + \sum_{\langle i,j \rangle} J_{ij}^{\mathrm{eff}}$$
 Sum over first neighbours of site i

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For a finite lattice:

$$\mu_i = \begin{cases} \mu_{\text{bulk}} = J'_{\text{eff}} + J_{\text{eff}} + U & \text{if} \quad 1 < i < 2M \\ \mu_{\text{edge}} = J'_{\text{eff}} + U & \text{if} \quad i \in \{1, 2M\} \end{cases}$$

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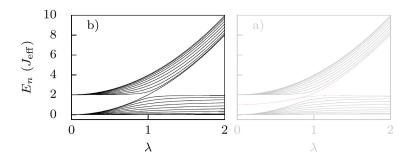
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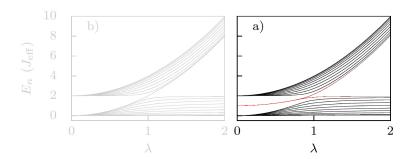
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This forbids the presence of edge states for doublons

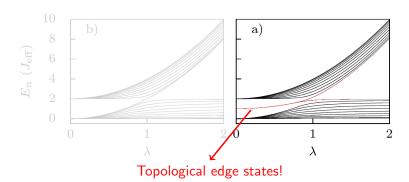
If we add a gate potential at the ends of the chain to compensate for the chemical potential difference, we recover the SSH model for doublons

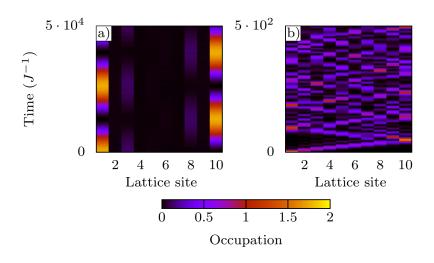


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$$U = 16J$$
; a) $\lambda = 0.5$, b) $\lambda = 1$

$$H(t) = H_{hopp} + H_U + E \cos(\omega t) \sum_{i=1}^{2M} x_i (n_{i\uparrow} + n_{i\downarrow})$$

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periodic driving

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For time-periodic Hamiltonians, H(t + T) = H(t), the fundamental solutions to the Schrödinger equation are of the form:

$$\psi(\mathbf{x}, t) = e^{-i\epsilon_n t} \phi_n(\mathbf{x}, t)$$

$$U(t_2, t_1) = e^{-iK(t_2)} e^{-iH_{eff}(t_2 - t_1)} e^{iK(t_1)}$$

$$H_{eff} = H^{[0]} + \frac{1}{\omega} H^{[1]} + \frac{1}{\omega^2} H^{[2]} + \cdots$$
(1)

M. Bukov, L. D'Alessio & A. Polkovnikov, Adv. Phys. 2015, Vol. 64, No. 2, 139-226

$$H(t) = H_{hopp} + H_U + E \cos(\omega t) \sum_{i=1}^{2NI} x_i (n_{i\uparrow} + n_{i\downarrow})$$

– Different effective Hamiltonians are obtained in the two regimes: $U\gg\omega$ and $\omega\gg U$

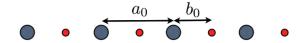
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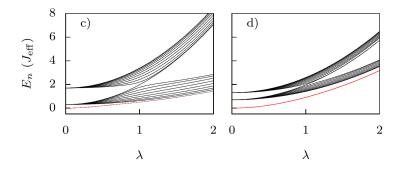
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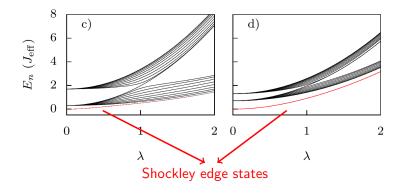
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– The ac field allows to tune the ratio between hopping and $\Delta\mu$ \Longrightarrow It is able to induce Shockley-like edge states

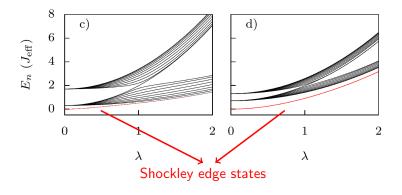
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These edge states also form a non-local two level system and allow the long-range transfer

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$$\begin{array}{c} 20 \\ 15 \end{array}$$

A. Gómez-León & G. Platero PRL 110, 200403 (2013)



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 - Driving the system with an AC field \rightarrow Shockley transfer
 - Combination of both \rightarrow AC induced topological transfer

Thank you!